

CMB power spectrum from cosmic strings

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N. Bevis, MH, M. Kunz^a astro-ph/0403029

N. Bevis, MH, M. Kunz, A. Liddle, J. Urrutilla in progress.

^aU. de Genève



University of Sussex

Introduction

String defects^a may be formed at phase transitions in the early Universe

- Thermal^b
- End of hybrid inflation^c
- $(D-\bar{D})$ -brane collisions^d

^aHindmarsh & Kibble (1994); Vilenkin & Shellard (1994); Kibble (2004)

^bKibble (1976); Zurek (1996); Rajantie (2002)

^cYokoyama (1989); Kofman, Linde, Starobinski (1996)

^dJones, Stoica, Tye (2002); Sarangi & Tye (2003); Copeland, Myers, Polchinski (2003)

... introduction

Observational consequences:

- Gravitational perturbations (scalar, vector & tensor)^a
- Baryon asymmetry^b
- Cosmic rays^c

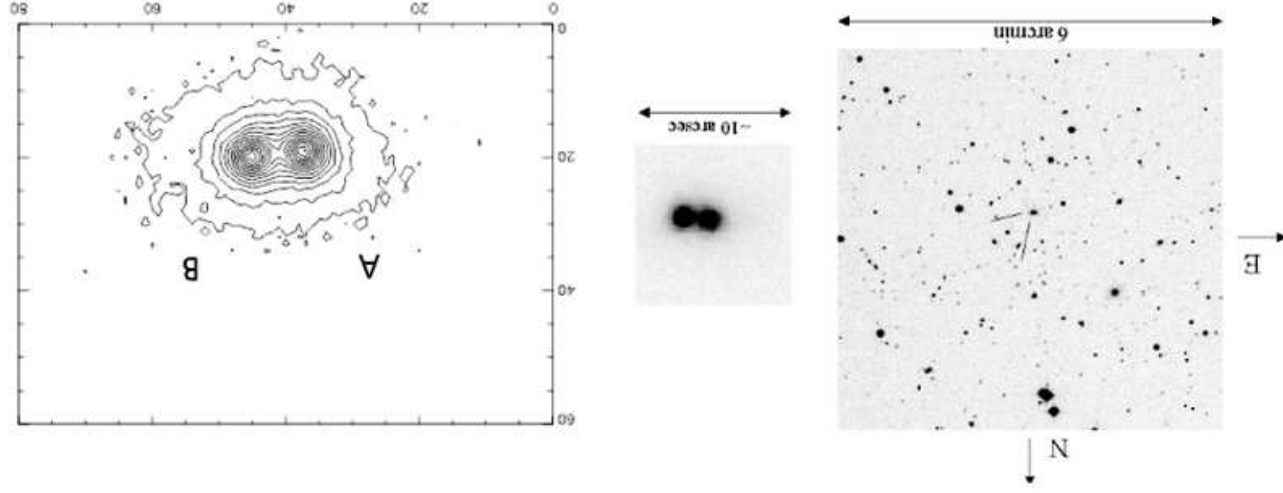
Network dynamics not well understood:
 Energy lost to (loops $\rightarrow$ gravitational waves) OR (classical radiation $\rightarrow$ particles)?
 Significant uncertainty in quantitative calculations (not often acknowledged)

^aZel'dovich (1980); Vilenkin (1981); Kaiser & Stebbins (1984); ...
^bBhattacharjee, Kibble, Turok (1982); Brandenburger, Davis, Hindmarsh (1991); Brandenburger, David, Trodden (1994); Sahu, Bhattacharjee, Yajnik (2004)
^cBhattacharjee (1990); Sigl (1996); Protheroe (1996); Berezhinski (1997); Vincent, M.H., Sakellariadou (1998)

New strings for old

Recent revival of interest:

- Cosmic F - and D -strings: M-theory allows many types of string network ^a
- CSL1: ^b



^aJones, Stoica, Tye (2002); Sarangi & Tye (2003); Copeland, Myers, Polchinski (2003)
^bSazhin et al (2003)

Topological defects in cosmology

- Global “defects”:^a
 - Spontaneous breaking of global symmetry
 - Low-energy dynamics: non-linear sigma model
 - Gravitational perturbations sourced by Goldstone modes
- Cosmic strings^b

- Spontaneous breaking of global or gauge symmetry
- Topology of vacuum manifold allows string-like solitons in 3D
- e.g. Abelian Higgs model (gauge cosmic strings)

^aDürrer, Kunz, Melchiorri *arXiv:astro-ph/0110348*

^bMH & Kibble *[arXiv:hep-ph/9411342]; Kibble arXiv:astro-ph/0410073; Polchinski arXiv:hep-th/0412244*

Global defects: temperature power spectrum

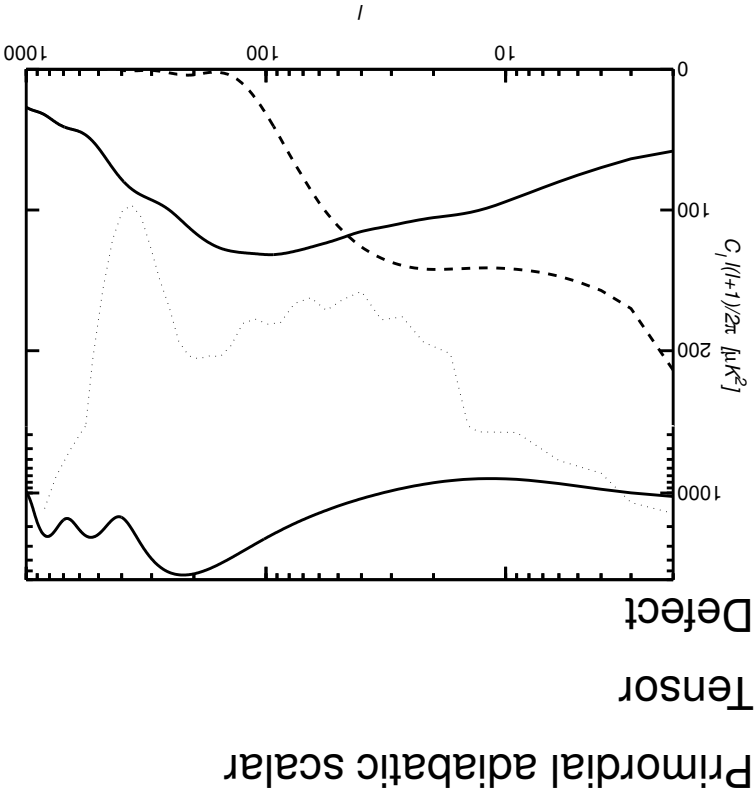
O(4) non-linear sigma model^a

Normalisation:

$C_{\ell}^{(\text{NLST})} = 0.13 C_{\ell}^{(\text{ad,scal.})}$
 $C_{\ell}^{(\text{tensor})} = 0.19 C_{\ell}^{(\text{scal.})}$.

Cosmological parameters:

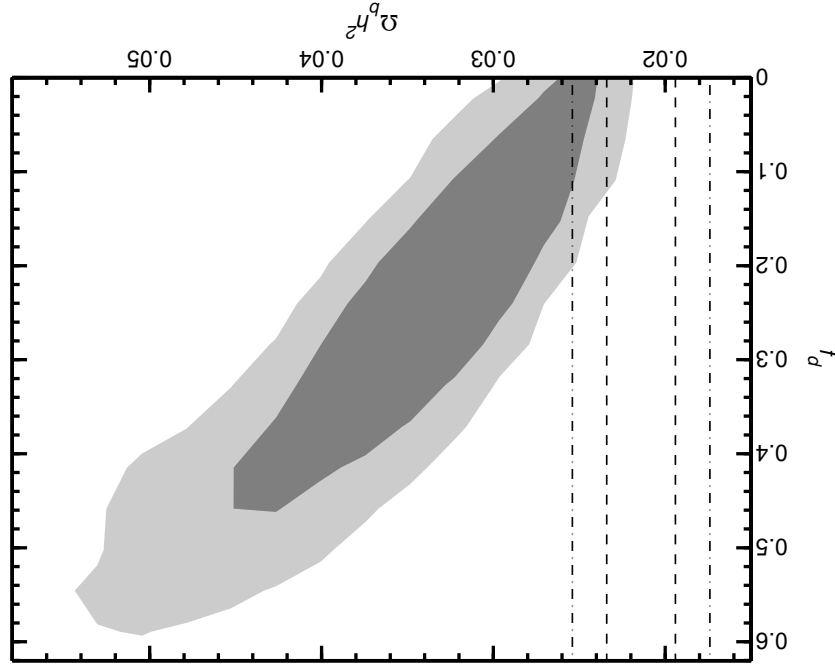
$\Omega = 1$
 $h = 0.70$
 $\Omega_b h^2 = 0.022$
 $\Omega_m = 0.30$
 $\tau = 0.10$ or $z_r = 13$



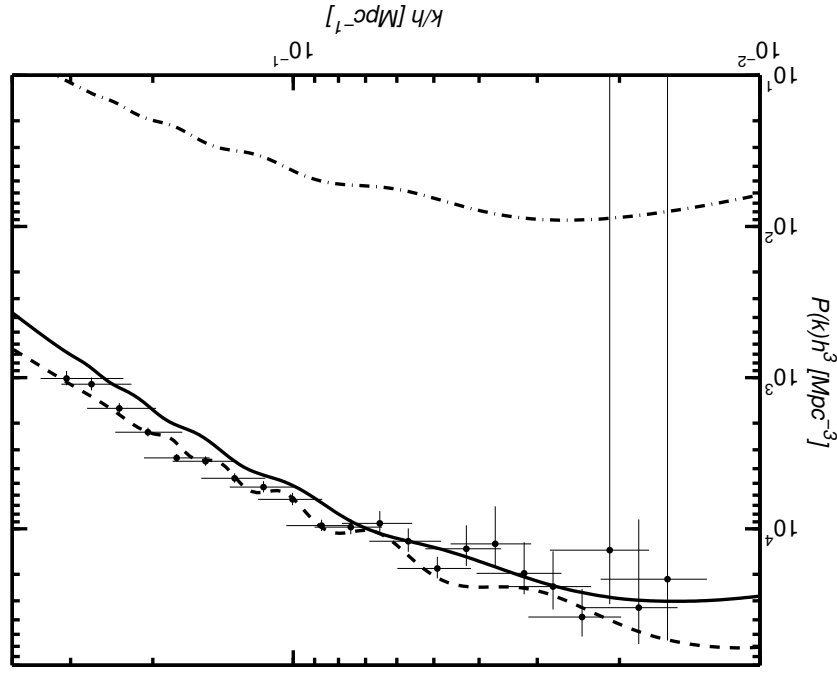
^aBevis, Hindmarsh, Kunz [arXiv:astro-ph/0403029]

Lifting the $\Omega_b h^2$ - f_d degeneracy with Big Bang Nucleosynthesis

- Marginalised likelihood from WMAP
- 68% and 95% confidence levels.
- Vertical lines: 68% and 95% confidence limits on $\Omega_b h^2$ (Kirkman *et al*)



SDSS and global defects: matter doesn't matter



Solid: Adiabatic HZ scalar
 Dashed: Scalar + $f_d = 0.4$
 Dot-dashed: defect contribution

Global defects: headline result

Global defects contribute less than 13%^a at $\ell = 9$ (95% confidence).

^a(theoretical uncertainty 2%)

Calculating perturbations from defects: UETC method

$h_a(\tau, k)$: linear perturbation (metric, matter, temperature ...)

$S_a(\tau, k)$: source (energy-momentum, separately conserved)

$D_{\alpha\beta}(\tau, k)$: time dependent differential operator

Perturbation equation: $D_{\alpha\beta}(\tau, k)h_\beta(\tau, k) = S_a(\tau, k)$

Power spectrum:^a $\langle |h_a(\tau_0, k)|^2 \rangle = \int \int \mathcal{D}^{-1} \mathcal{D}^{-1} \langle S_a(\tau, k) S_a^*(\tau', k) \rangle$

Need unequal-time correlators (UETCs) of source or energy-momentum tensor

$$C^{\mu\nu\rho\lambda}(k, \tau, \tau') = \langle T^{\mu\nu}(k, \tau) T^{\rho\lambda}_*(k, \tau') \rangle$$

^a Pen, Seljak, Turok (1997); Dürer, Kunz, Melchiorri (1998, 2002)

... UETC method

Isotropy + EM conservation + parity: $C^{\mu\nu\rho\lambda}$: 3 scalar, 1 vector, 1 tensor

(S,V,T) correlators can be diagonalised: $C(k_\tau, k_{\tau'}) = \sum_n \lambda_n v_n(k_\tau) v_n^*(k_{\tau'})$

Eigenvectors $v_n(k_\tau)$ as sources: $h_n^a(\tau, k) = \int_\tau^{\tau_i} \mathcal{D}_{-1}^{\alpha\beta}(\tau, \tau', k) v_n(\tau', k)$

Reconstruct complete power spectrum:

$$C_\ell = \sum_n \lambda_{(S)}^n C_{(S)\ell}^n + \sum_n \lambda_{(V)}^n C_{(V)\ell}^n + \sum_n \lambda_{(T)}^n C_{(T)\ell}^n$$

UTC method for power spectrum: summary

Calculate UTCs



Diagonalise UTCs



Solve perturbation equations with eigenfunctions as sources

Square (e.g.) $(\Delta T)^\ell (S, V, T)^n$ and sum

Energy momentum tensor decomposition

Scalar (S), Vector (V) and Tensor (T) under 3D rotation group^a

$$\begin{aligned}
 T_{(S)}^{00} = v^2 f_\rho & \quad T_{(S)}^{j0} = v^2 k_j f_v & T_{(S)}^{jl} = v^2 \left[(f_p + \frac{1}{3} k^2 f_\pi) \delta_{jl} - k_j k_l f_\pi \right] \\
 T_{(V)}^{j0} = v^2 w_j & \quad T_{(V)}^{jl} = \frac{1}{2} v^2 \left(k_j w_l^{(\pi)} + k_l w_j^{(\pi)} \right) \\
 T_{(T)}^{jl} & = v^2 T_{ij}^{(\pi)}
 \end{aligned}$$

v is v.e.v. of scalar field: $V(\phi) = \frac{1}{2} \lambda (|\phi|^2 - v^2)^2$

^aNotation of Dürer, Kunz, Melchiorri [arXiv:astro-ph/0110348]

Linearised Einstein equations for defect “seeds”

Gauge invariant formalism^a

Scalar: $\Phi^{(s)}$, $\Psi^{(s)}$

Vector: $\Sigma^i_{(s)}$

Tensor: $H^{ij}_{(s)}$

$$\begin{aligned} k_2 \Phi^{(s)} &= \epsilon(f_\rho + 3 \frac{a}{f_v}) \\ -2 \epsilon f_\pi &= \Phi^{(s)} + \Psi^{(s)} \\ -k_2 \Sigma^i_{(s)} &= 4 \epsilon w^i_{(v)} \\ H^{ij}_{(s)} + 2 \frac{a}{f_v} H^{ij}_{(s)} + k_2 H^{ij}_{(s)} &= 2 \epsilon \tau^{ij}_{(v)} \end{aligned}$$

$\epsilon = 4\pi G v^2$, where v is v.e.v. of scalar field

^aNotation of Dürer, Kunz, Melchiorri [arXiv:astro-ph/0110348]

Do not have to trace entire history of network

Defect networks exhibit **scaling**

Scaling: correlators depend only on time τ and wavenumber k

$$\begin{aligned} \langle \Phi_s(\mathbf{k}, \tau) \Phi_s^*(\mathbf{k}', \tau) \rangle &= \frac{1}{k^4 \sqrt{\tau \tau'}} C_{11}(x, x') \\ \langle \Phi_s(\mathbf{k}, \tau) \Psi_s^*(\mathbf{k}', \tau) \rangle &= \frac{1}{k^4 \sqrt{\tau \tau'}} C_{12}(x, x') \\ \langle \Psi_s(\mathbf{k}, \tau) \Psi_s^*(\mathbf{k}', \tau) \rangle &= \frac{1}{k^4 \sqrt{\tau \tau'}} C_{22}(x, x') \end{aligned}$$

NB scaling fails during a change in expansion rate

^aNotation of Dürer, Kunz, Melchiorri [arXiv:astro-ph/0110348]

Unequal time correlators: scaling

CMB power spectrum from (gauge) cosmic strings: previous work

Authors	String model	COBE G_{μ}
Perivolaropoulos (1995)	Moving segment model	$\sim 1.6 \times 10^{-6}$
Allen et al (1997)	Ideal strings (FRW)	$1.05^{+0.35}_{-0.20} \times 10^{-6}$
Albrecht, Battye, Robinson (1997)	Moving segment model	$\sim 2 \times 10^{-6}$
Contaldi, Hindmarsh, Magueijo (1998)	Ideal strings (Minkowski)	—
Landraiu & Shellard (2004)	Ideal strings (FRW)	$\sim 0.7 \times 10^{-6}$
Wyman, Pogosian, Wasserman (2005)	Moving segment model	$\sim 2 \times 10^{-6}$

MSM: moving segments of increasing length with random velocities^a

Ideal Strings: 1D objects, Nambu-Goto action

^aVincent, Hindmarsh, Sakellariadou (1997); Albrecht, Battye, Robinson (1997)

Approximations

Quantum field theory

↑ (Large occupation number, bosons)

Classical field theory

↑ (Smooth, low curvature string configurations)

“Ideal” (Nambu-Goto) strings

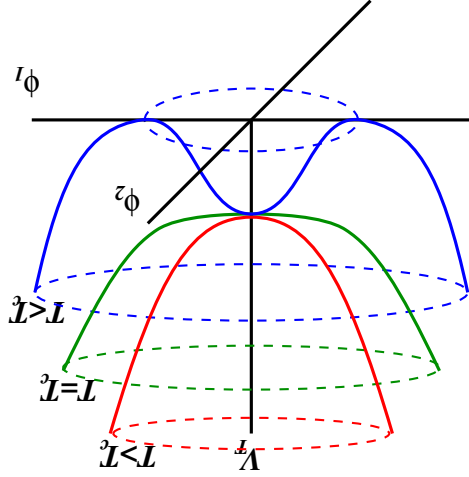
↑ (Phenomenological)

Moving segment model

Abelian Higgs model

$$S = - \int d^4x \sqrt{-g} \left(g_{\mu\nu} D_\mu \phi^* D_\nu \phi + V(\phi) + \frac{1}{4e^2} g_{\mu\rho} g_{\nu\sigma} F^{\mu\nu} F^{\rho\sigma} \right),$$

Complex scalar field $\phi(\mathbf{x}, t)$, vector field $A_\mu(\mathbf{x}, t)$
 Covariant derivative $D_\mu = \partial_\mu - ieA_\mu$.
 Potential $V(\phi) = \frac{1}{2}\lambda(|\phi|^2 - v^2)^2$.
 Metric $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$



Temporal gauge ($A_0 = 0$) field equations (index raised with Minkowski metric).

$$\ddot{\phi} + 2\frac{\dot{a}}{a}\dot{\phi} - D^2_2\phi + \lambda a^2(|\phi|^2 - v^2)\phi = 0,$$

$$\partial_\mu \left(\frac{1}{e^2} F^{\mu\nu} \right) - ia^2(\phi^* D_\nu \phi - D_\nu \phi^* \phi) = 0,$$

Abelian Higgs model: shrinking string problem

Comoving width of string shrinks as a^{-1} ($a \sim \tau^{-2}$ in matter era)

Modify field equations^a

$$\ddot{\phi} + 2\frac{\dot{a}}{a}\dot{\phi} - D^2\phi + \lambda a^{2s}(|\phi|^2 - v^2)\phi = 0,$$

$$\partial_\mu \left(a^{2(1-s)} \frac{e^2}{F^{\mu\nu}} - ia^2(\phi^* D_\nu \phi - D_\nu \phi^* \phi) \right) = 0,$$

Physical width of string “fattens” if $s > 1$

Preserves Gauss’s Law (current conservation) but violates EM conservation

Take $s = 0.3$ (check with $s = 0$)

^aPress, Ryden, Spergel (1989); Bevis et al in prep.

Parallel simulations of field theories: LATfield

- C++ library of objects for classical lattice simulation^a
- Inspired by MDP/FermiQCD^b
- Objects:

Lattice: Takes care of boundary conditions and domain decomposition

Field: Template - can have real, complex, user-defined object.

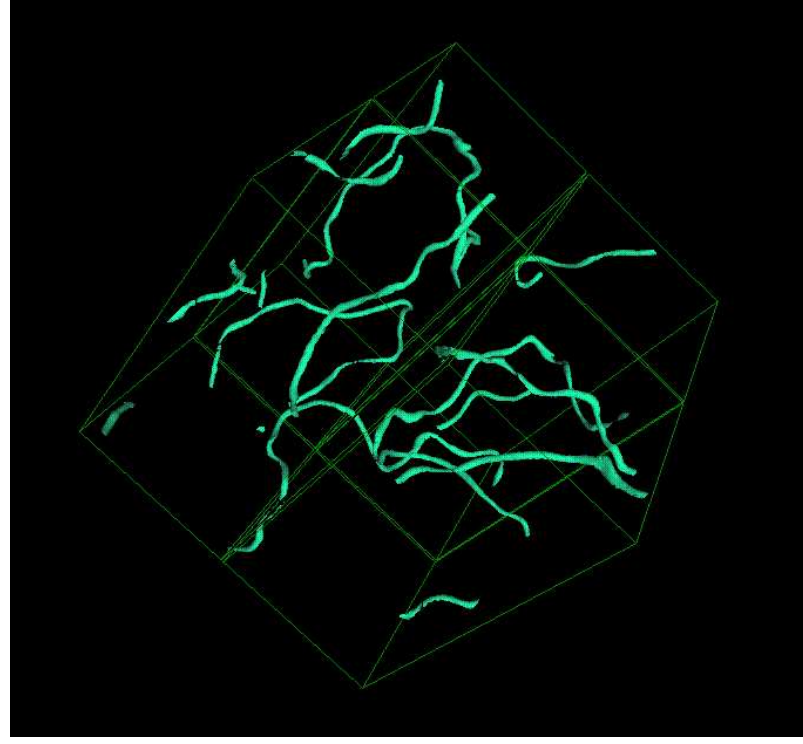
Site: Accesses elements of field

- Parallelisation: compiler switch – DPARALLEL_MP I

^a Bevis & Hindmarsh, in preparation

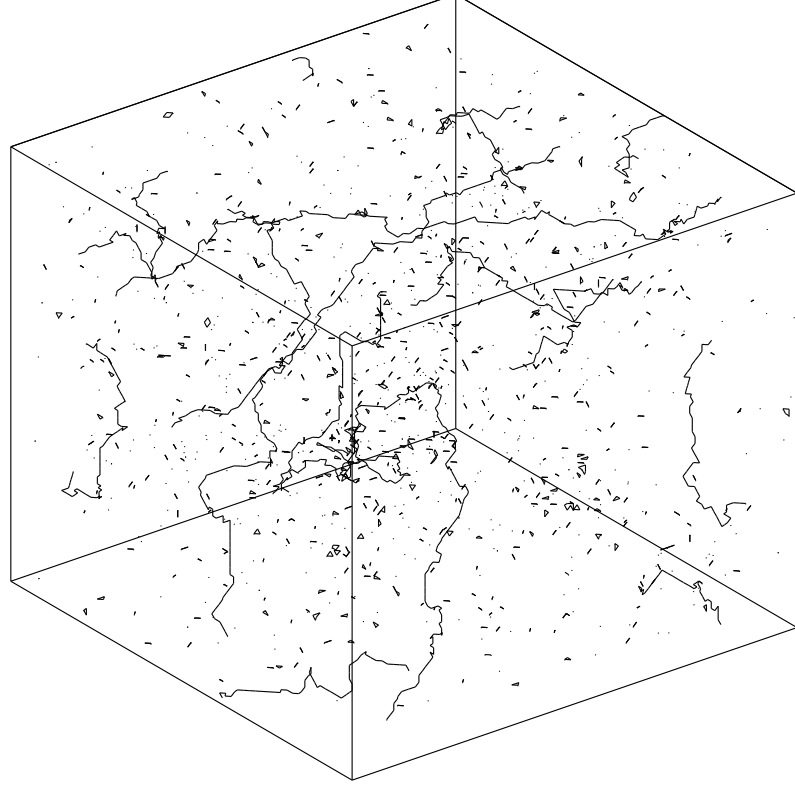
^b Massimo di Pierro <http://www.fermiqcd.net/>

Strings in 3D (Minkowski space)



Classical Abelian Higgs model

Isosurfaces of constant ϕ
 ($\phi = 0$ at string centre).



Nambu-Goto strings

Abelian Higgs model simulations 1: string length scale

Total length of string: L

Scaling: $L/V \propto \tau^{-2}$

Network scale: $\xi = \sqrt{(V/L)} \propto \tau$

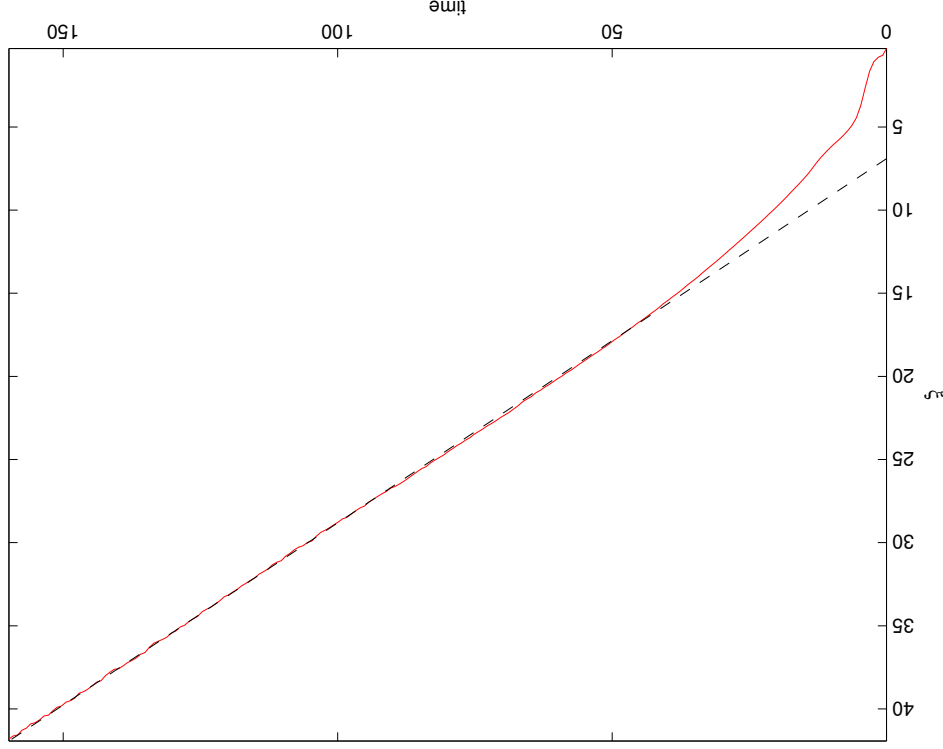
Lattice spacing: $\Delta x = 0.5$

Volume 512^3

Coupling constants: $\lambda = e = 1$

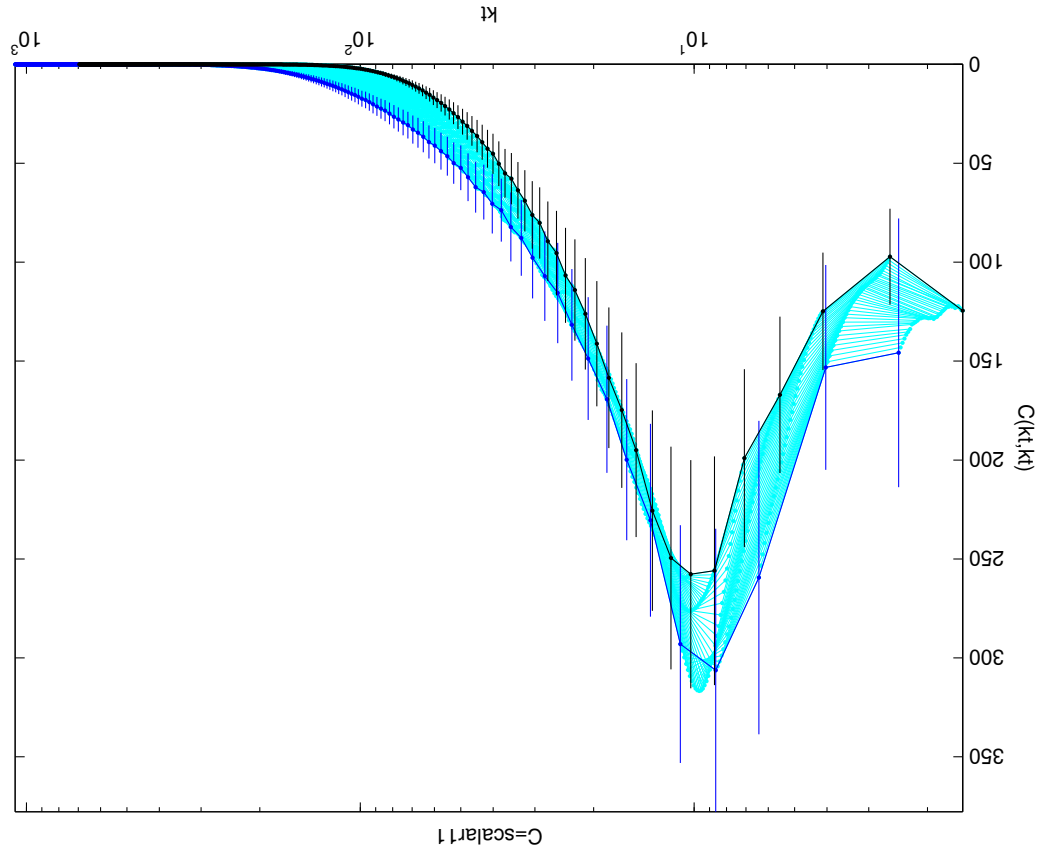
Matter era, “fattening” string algorithm.

Note $\xi = x(\tau + \tau_0)$



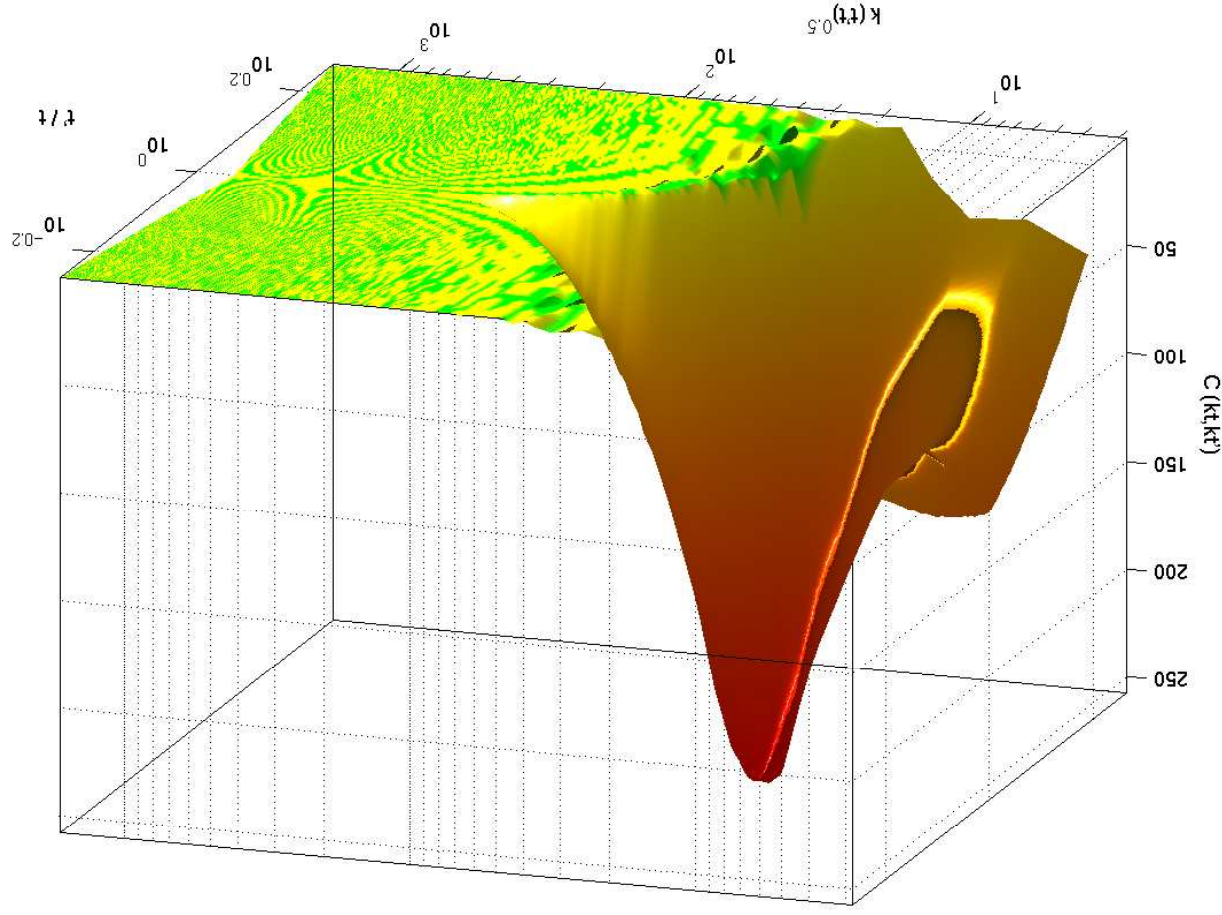
$$C_{11}(k(\tau + \tau_0), k(\tau + \tau_0)) = k^4(\tau + \tau_0) \langle \Phi_s(\mathbf{k}, \tau) \Phi_s^*(\mathbf{k}, \tau) \rangle$$

$\Phi\Phi$ equal time correlator: apply time offset τ_0 , start at $\tau = 64$

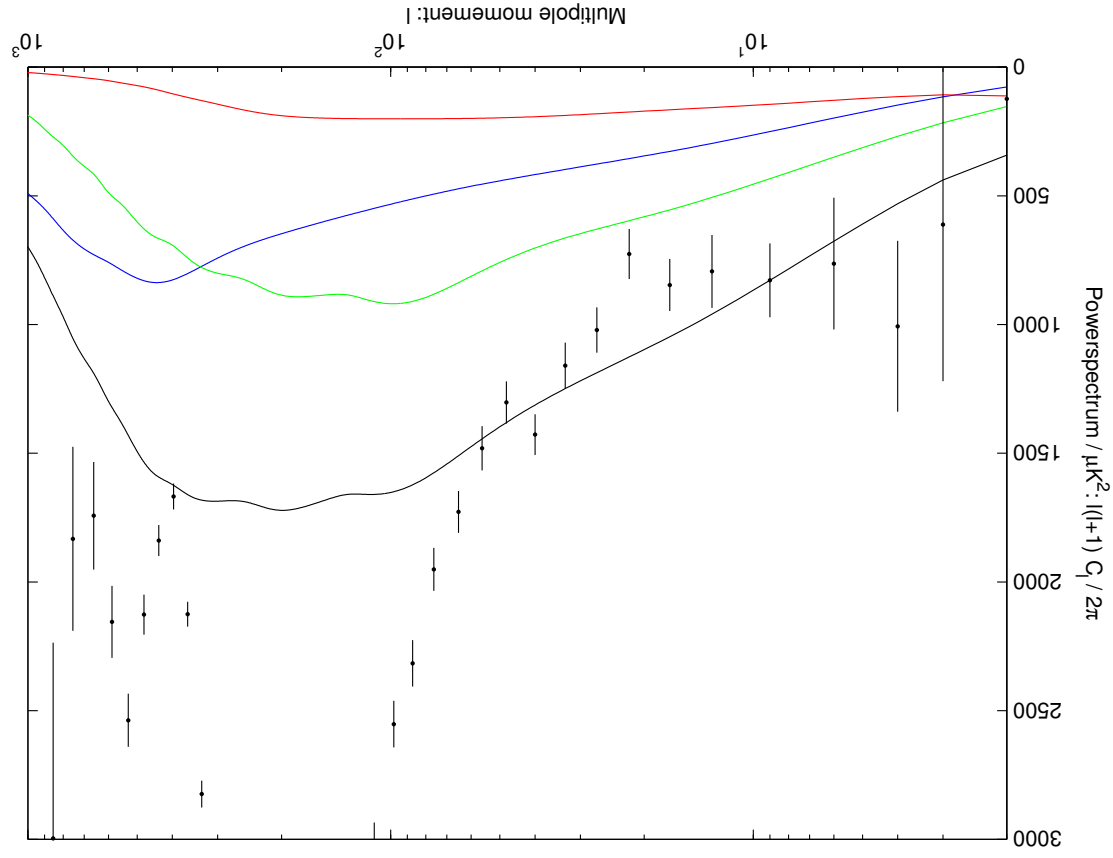


Simulations 2: Φ unequal time correlator

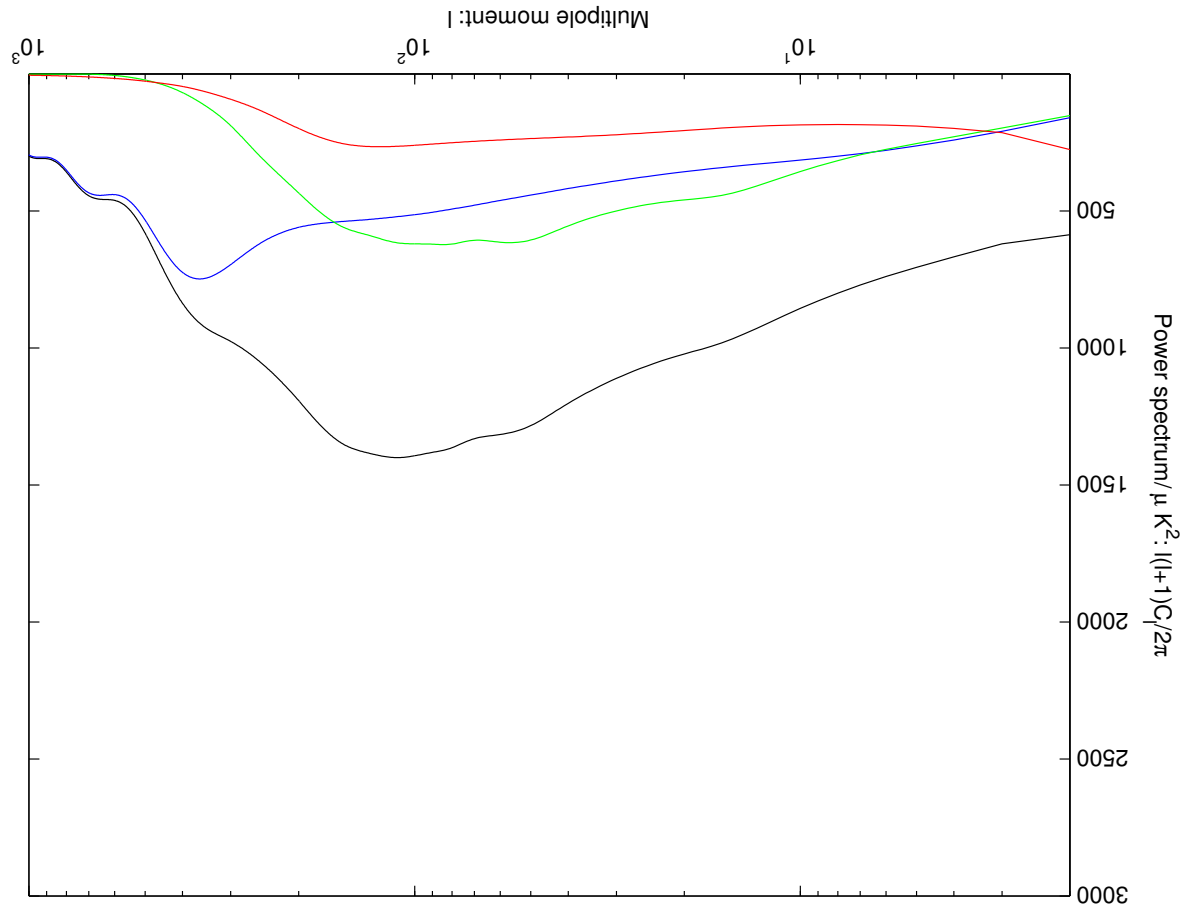
Plot against $k\sqrt{\tau\tau'}$, τ'/τ instead of $k\tau$, $k\tau'$.



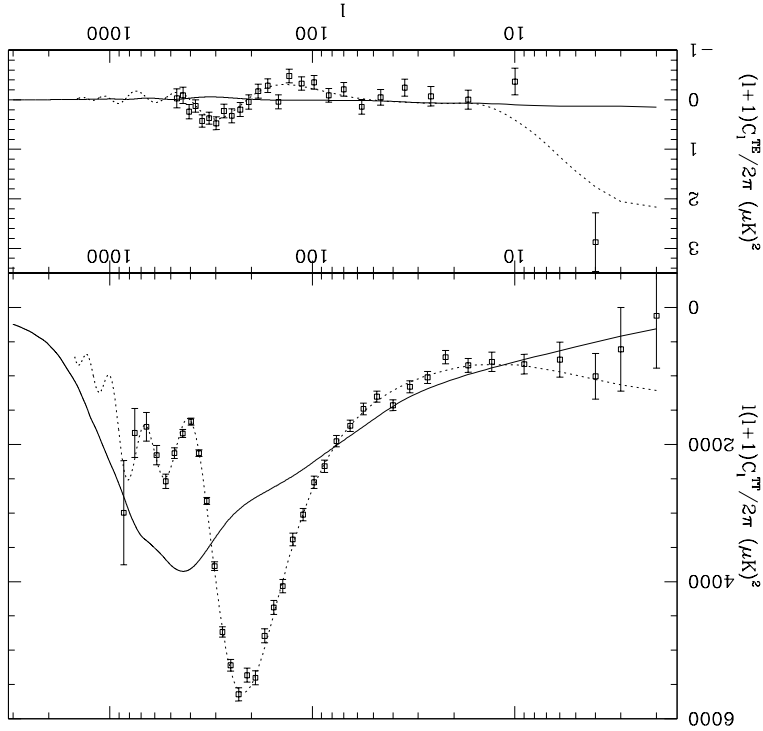
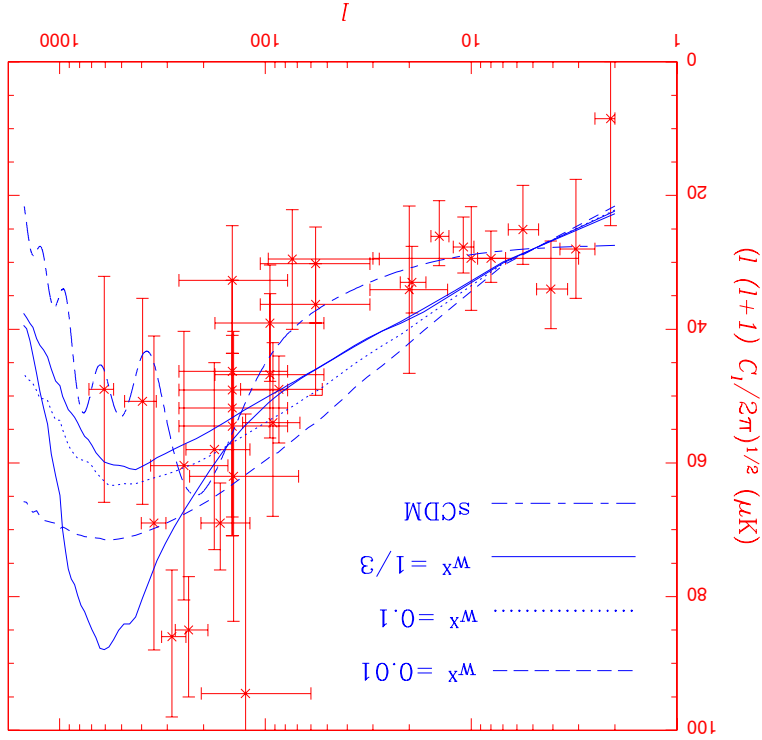
WMAP-normalised C_ℓ s from strings PRELIMINARY!



WMAP-normalised C_ℓ s from O(4) NLSM



Previous calculations of cosmic string C_ℓ s



Contaldi, Hindmarsh, Magueijo (1998) Wyman, Pogosian, Wasserman (2005) Minkowski Nambu-Goto Moving segment model

Conclusions

- Cosmic strings may be detected as a subdominant contribution to CMB signal
- **NEW:** field theory simulations:
 - less modelling, more physics
 - include contributions to EM tensor from decay products
- Technology spin-offs:
 - parallel N -dimensional field theory simulations: **LATfield**
 - codes for UETCs and sourced perturbations (modified CMBeasy)
- Preliminary results:
 - Larger vector contribution
 - broader, lower, peak
- To be done:
 - MCMC parameter estimation to bound symmetry-breaking scale
 - Statistical, finite size, and model error estimates