

MeV-scale Reheating and Neutrino Oscillations

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@COSMO-05

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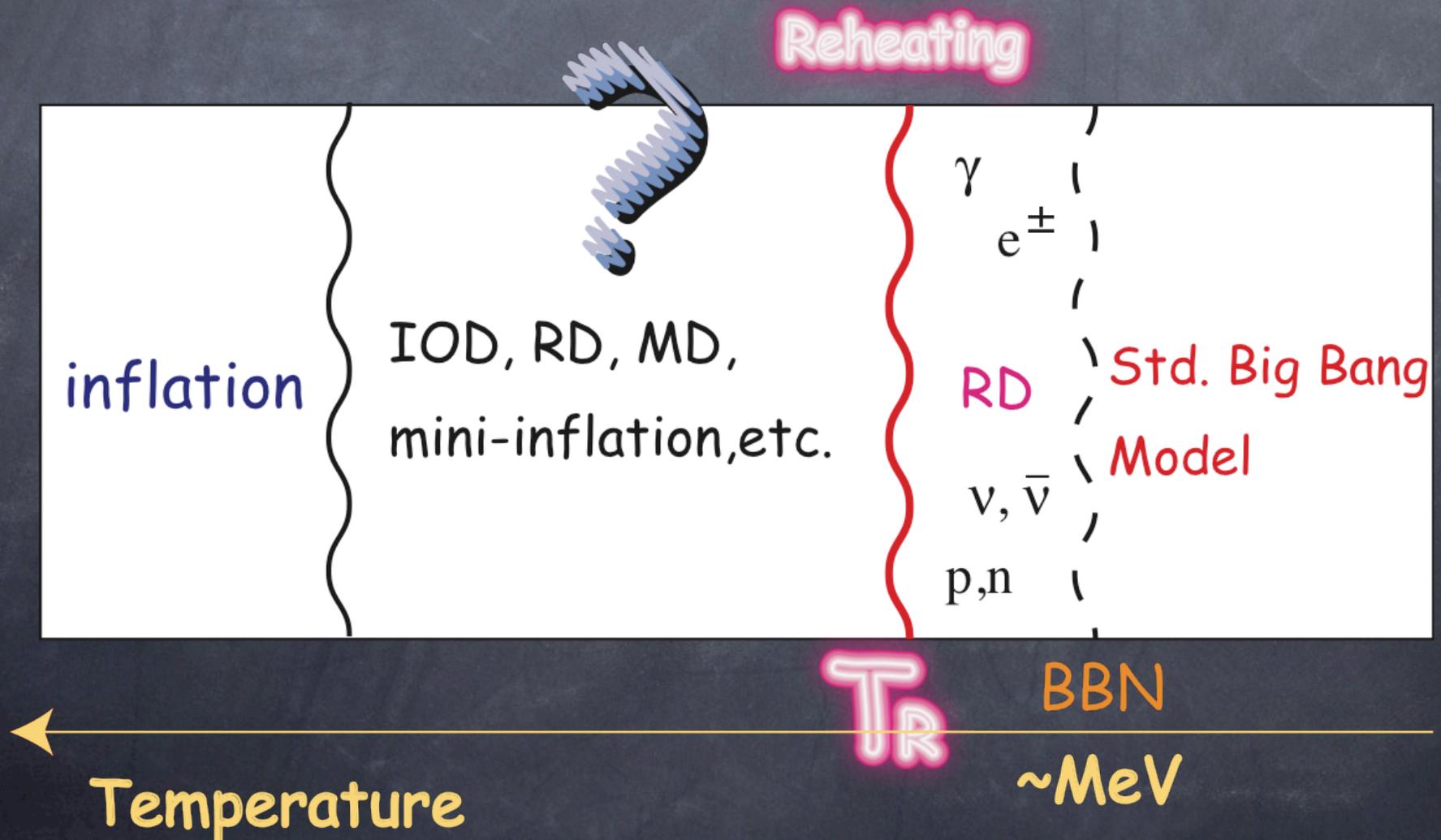
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K. Ichikawa, M. Kawasaki, F.T. astro-ph/0505395 (to be published in PRD)

I. Introduction

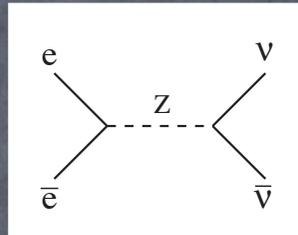
- What do we know about thermal history of the universe?



What if T_{RH} is around MeV?

- ✓ Neutrinos might not be fully thermalized.

[Kawasaki, Kohri, Sugiyama, '99, '00]



The predicted abundances of the light elements (especially ${}^4\text{He}$) are affected.

It has been widely believed that $N_\nu \curvearrowright Y_p \curvearrowright$

What we did:

- We investigated the MeV-scale reheating scenario, taking account of thermalization processes of neutrinos

+ neutrino oscillations

- ✓ We found that ${}^4\text{He}$ abundance **increases** as N_ν decreases due to flavor mixings!



(to be detailed later)

II. Role of neutrinos in BBN

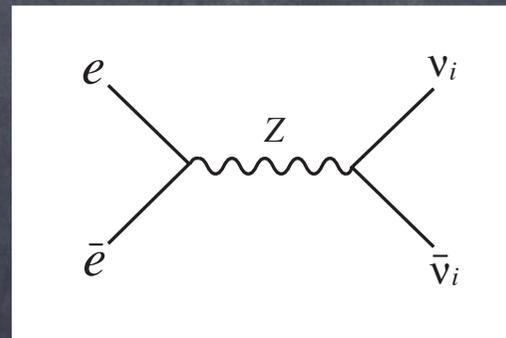
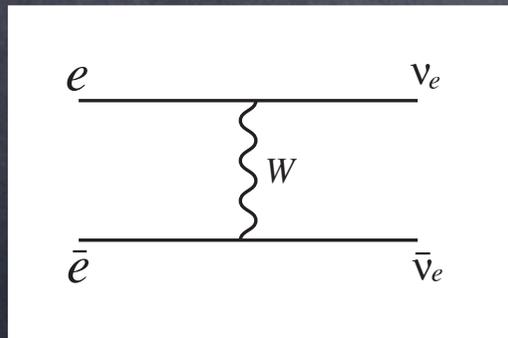
Reheating processes (Setup)

$$\phi \rightarrow \gamma, e^{\pm}$$

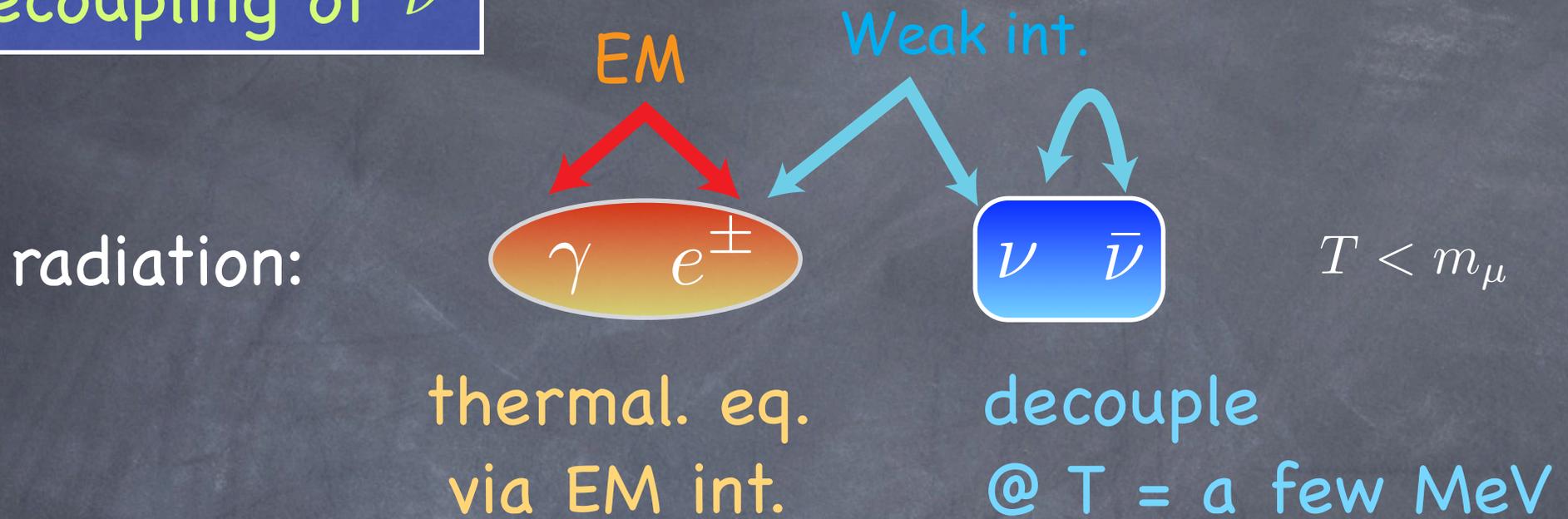
We define the reheating temperature by $\Gamma = 3H(T_R)$

$$\Gamma = 2.03 \left(\frac{T_R}{\text{MeV}} \right)^2 \text{sec}^{-1}.$$

- We assume that neutrinos are exclusively produced via $e^- + e^+ \leftrightarrow \nu + \bar{\nu}$



Decoupling of ν



In std. BBC,

$$T_d \sim 5\text{MeV} \quad \text{for } \nu_{\mu,\tau}$$

$$T_d \sim 3\text{MeV} \quad \text{for } \nu_e$$

[Hannestad & Madsen '95]

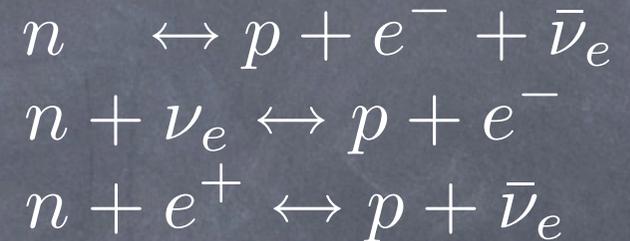
[Dolgov, Hansen & Semikoz, '97, '99]

ν_e is more likely to be produced at $T \sim \text{MeV}$.

${}^4\text{He}$ abundance

n-p transformation decouples when

$$\Gamma_{n \leftrightarrow p} = H$$



→ n/p ratio fixes (except for neutron free decay) at $T \sim 1 \text{ MeV}$.

$$\left(\frac{n}{p}\right)_{\text{EQ}} = \exp\left(-\frac{Q}{T}\right)$$

$$Q = m_n - m_p = 1.293 \text{ MeV}$$

→ Almost all neutrons are absorbed in ${}^4\text{He}$, soon after the deuterium bottleneck opens at $T_D \sim 0.08 \text{ MeV}$.

What if $N_\nu < 3$?

$$N_\nu \equiv \Sigma \rho_\nu / \rho_{\nu, \text{std.}}$$

✓ Hubble parameter decreases.

H



Y_p



✓ The n-p tranf. rate also decreases.

$\Gamma_{n \leftrightarrow p}$



Y_p



✓ D bottleneck opens later.

$H(T_D)^{-1}$



Y_p



III. Thermalization of Neutrinos

• Density matrices:

$$\begin{aligned}\langle a_j^\dagger(\mathbf{p}) a_i(\mathbf{p}') \rangle &= (2\pi)^3 \delta^{(3)}(\mathbf{p} - \mathbf{p}') [\rho_{\mathbf{p}}]_{ij}, \\ \langle b_i^\dagger(\mathbf{p}) b_j(\mathbf{p}') \rangle &= (2\pi)^3 \delta^{(3)}(\mathbf{p} - \mathbf{p}') [\bar{\rho}_{\mathbf{p}}]_{ij}.\end{aligned}$$

$$\{i, j\} = \{e, \mu, \tau\}$$

- diagonal components: dist. func. of ν_i
- off-diagonal ones: corr. between ν_i and ν_j .

QKEs for density matrix

$$\begin{aligned} \nu + e^\pm &\leftrightarrow \nu + e^\pm \\ \nu + \bar{\nu} &\leftrightarrow e^- + e^+ \end{aligned}$$

$$Ha \frac{d\rho_p}{da} = -i[\Omega(p), \rho_p] + I_{\text{coll}}(p).$$

refractive term
[flavor mixings]

collision term
[neutrino thermalization]

where

$$\Omega(p) \equiv \Omega_V(p) - \frac{8\sqrt{2}G_F p}{3m_W^2} E$$

$$\Omega_V(p) = \frac{1}{2p} U M^2 U^T, \quad M^2 \equiv \begin{pmatrix} m_1^2 & 0 \\ 0 & m_2^2 \end{pmatrix}, \quad U \equiv \begin{pmatrix} \cos \theta_{12} & \sin \theta_{12} \\ -\sin \theta_{12} & \cos \theta_{12} \end{pmatrix}.$$

$$E = \begin{pmatrix} \rho_e + \rho_{\bar{e}} & 0 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} (7/60)\pi^2 T^4 & 0 \\ 0 & 0 \end{pmatrix},$$

Approximate 2 flavor analysis
assuming $\theta_{13} \approx 0$

Parameters:

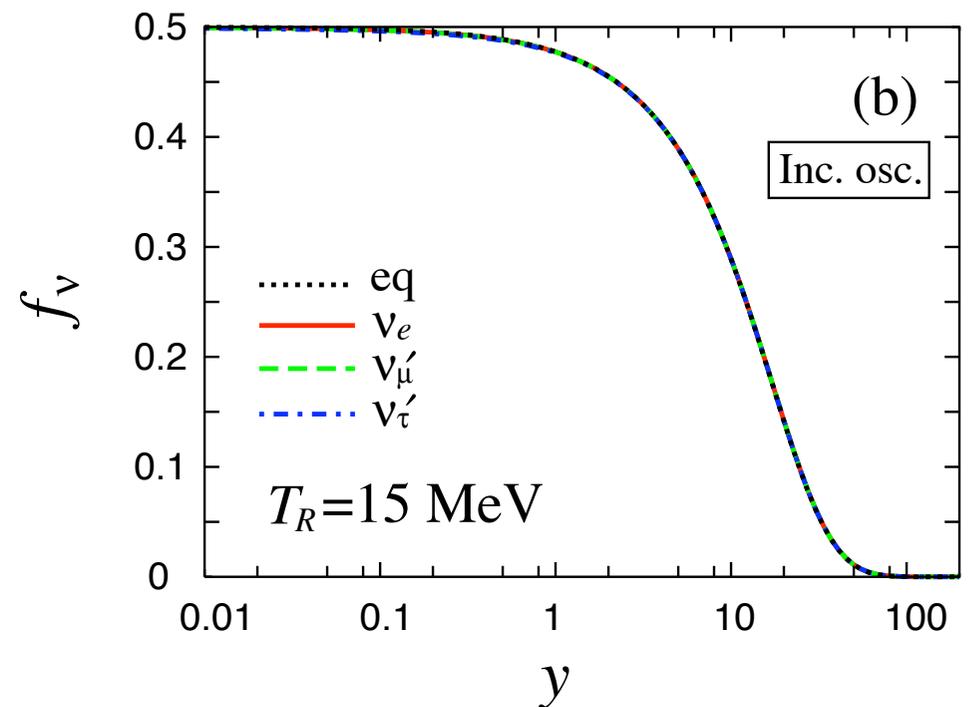
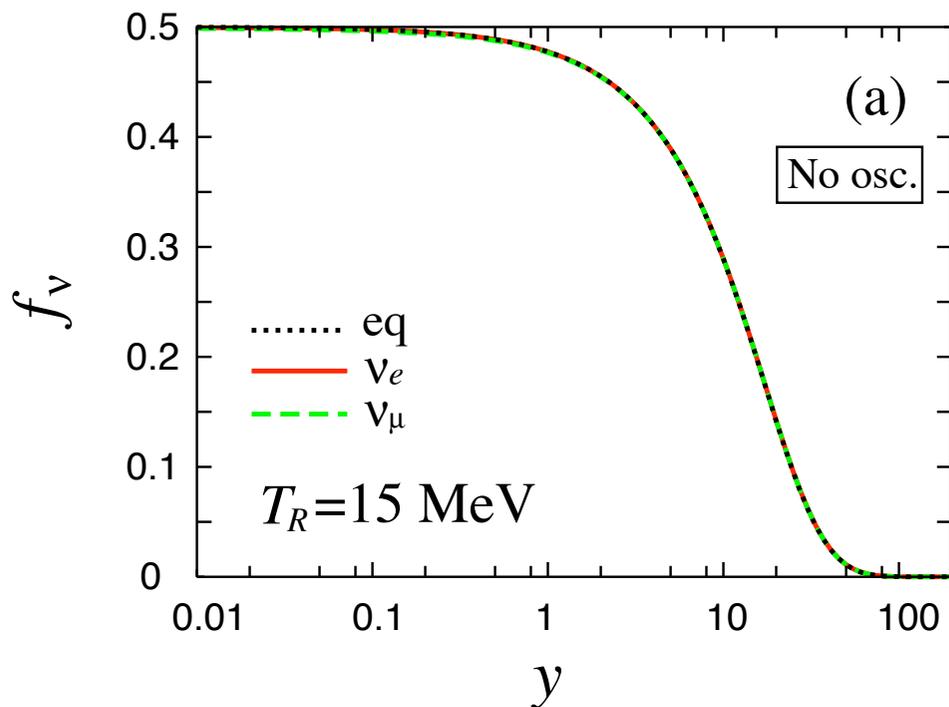
$$\sin^2 \theta_{12} = 0.315$$

$$m_2^2 - m_1^2 = 7.3 \times 10^{-5} \text{eV}^2$$

IV. Results & Discussions

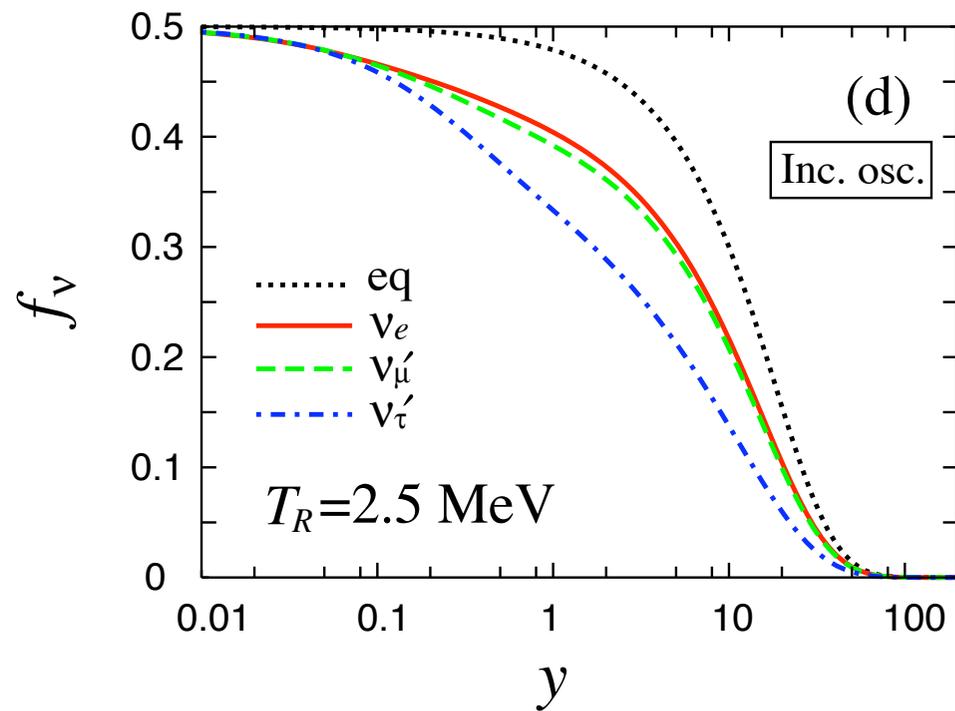
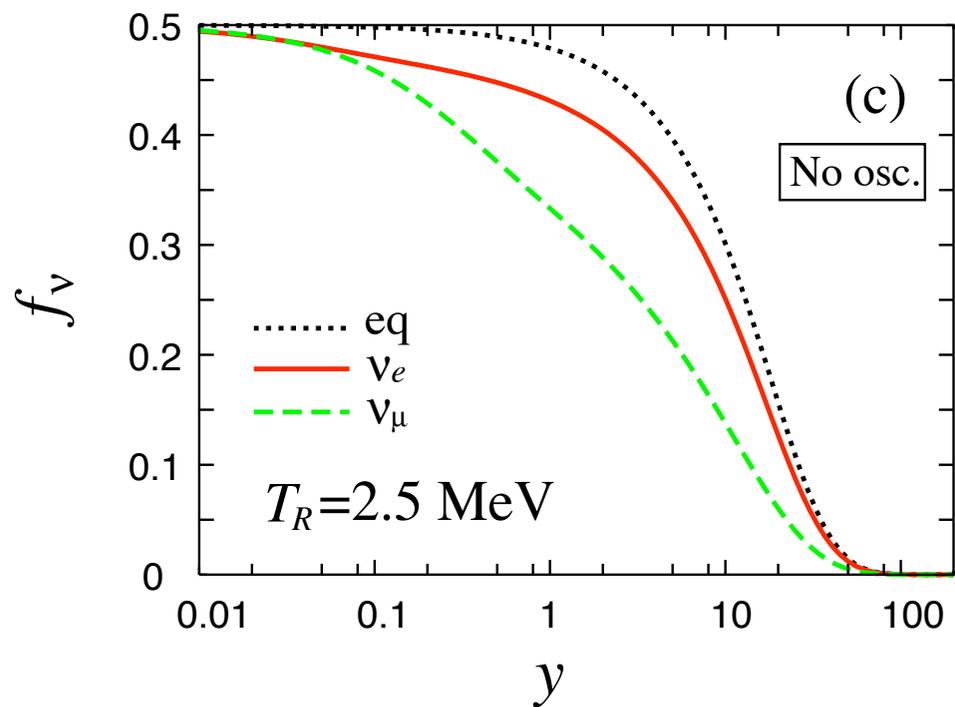
- The final dist. functions [$T_R=15\text{MeV}$]

Almost thermal distribution!



Momentum is normalized as $y = ap$

The final dist. functions [$T_R=2.5\text{MeV}$]

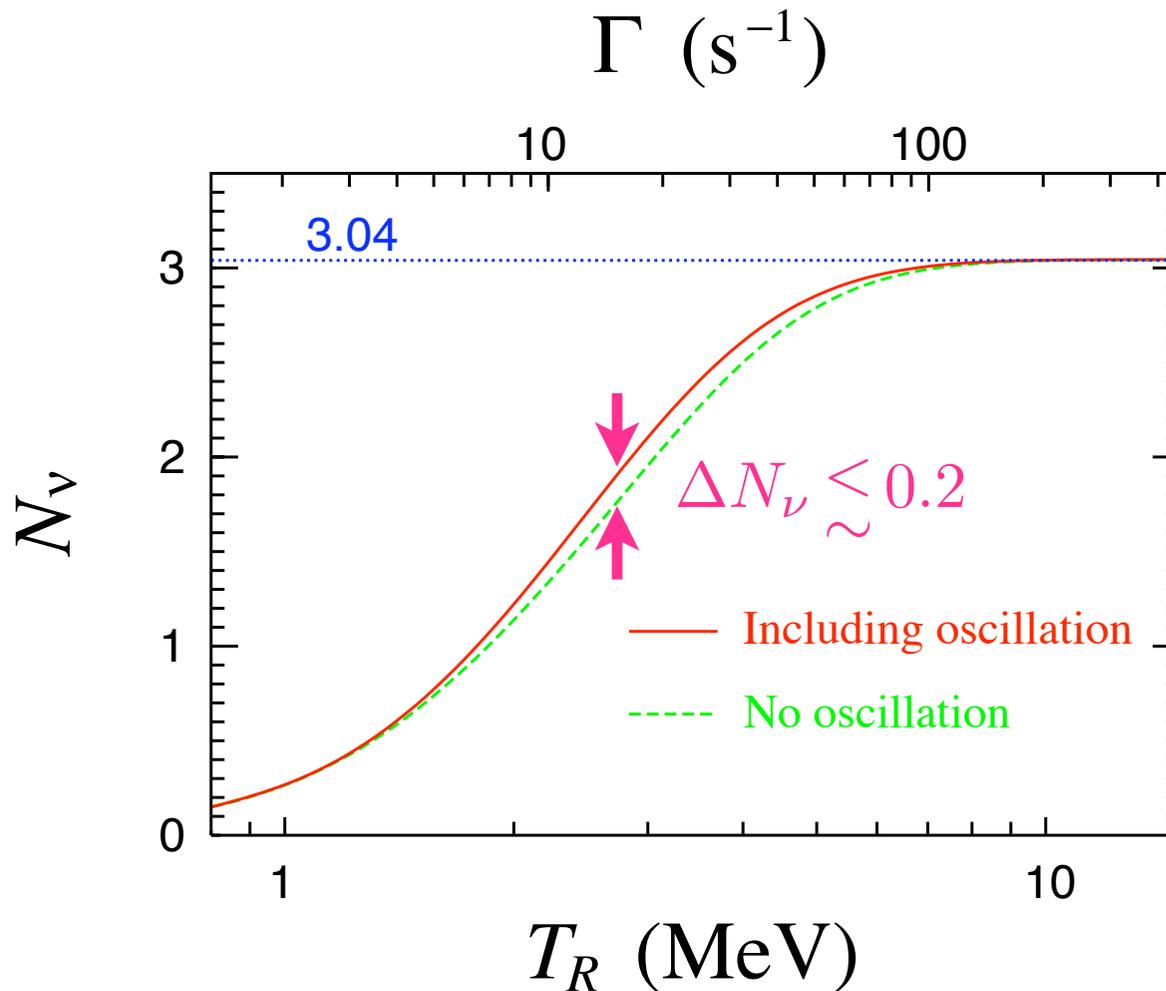


$$f_{\nu_e} > f_{\nu'_\mu} = f_{\nu'_\tau}$$

$$f_{\nu_e} \sim f_{\nu'_\mu} > f_{\nu'_\tau}$$

Due to flavor mixings, $f_{\nu'_\mu} \nearrow f_{\nu_e} \searrow \Gamma_{n \leftrightarrow p} \searrow$

Effective number of neutrinos



N_ν (i.e. H) does not change much!

Two competing effects:

(1) Slowing down of the expansion rate

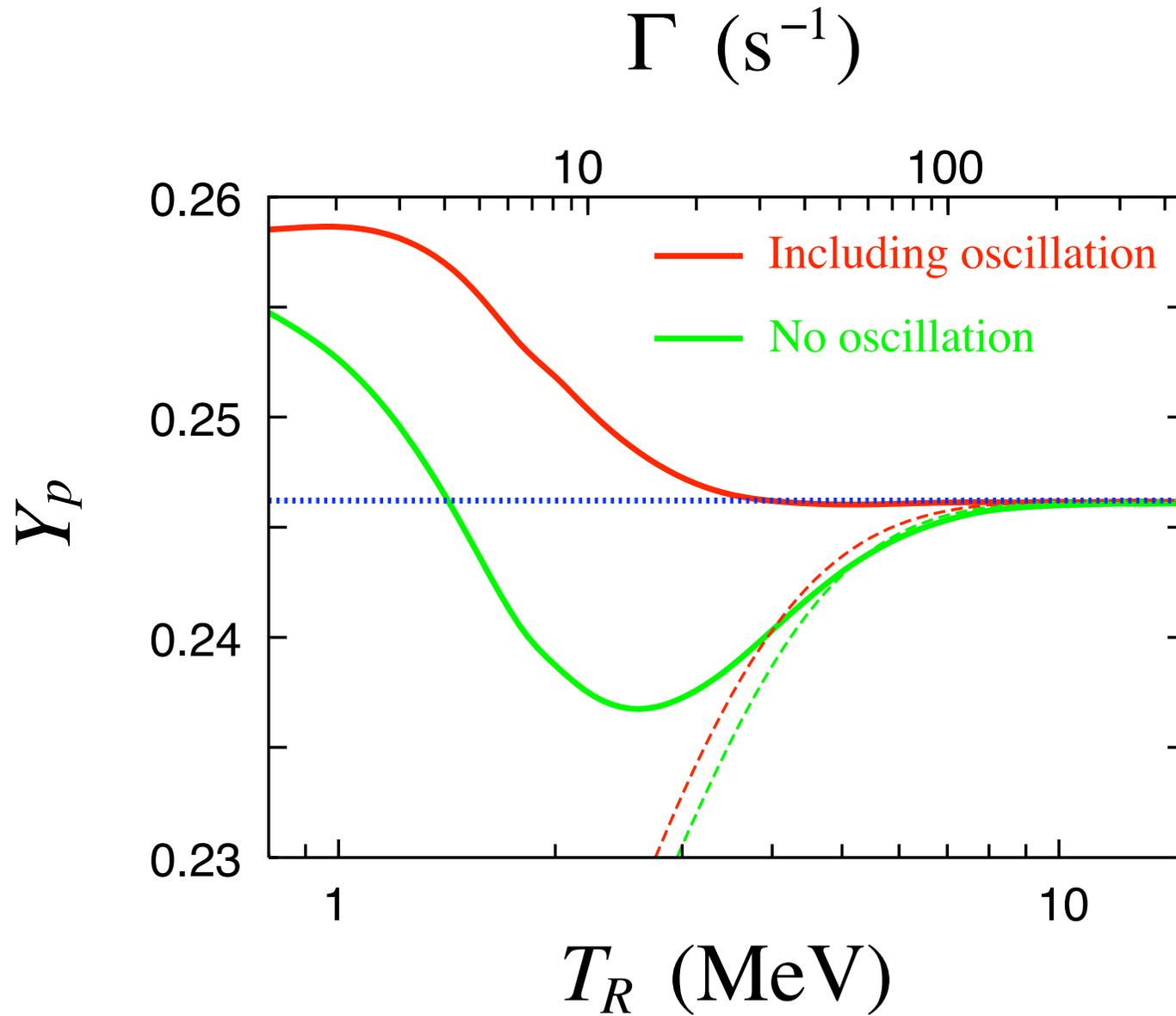


(2) Decrease in $\Gamma_{n \leftrightarrow p}$

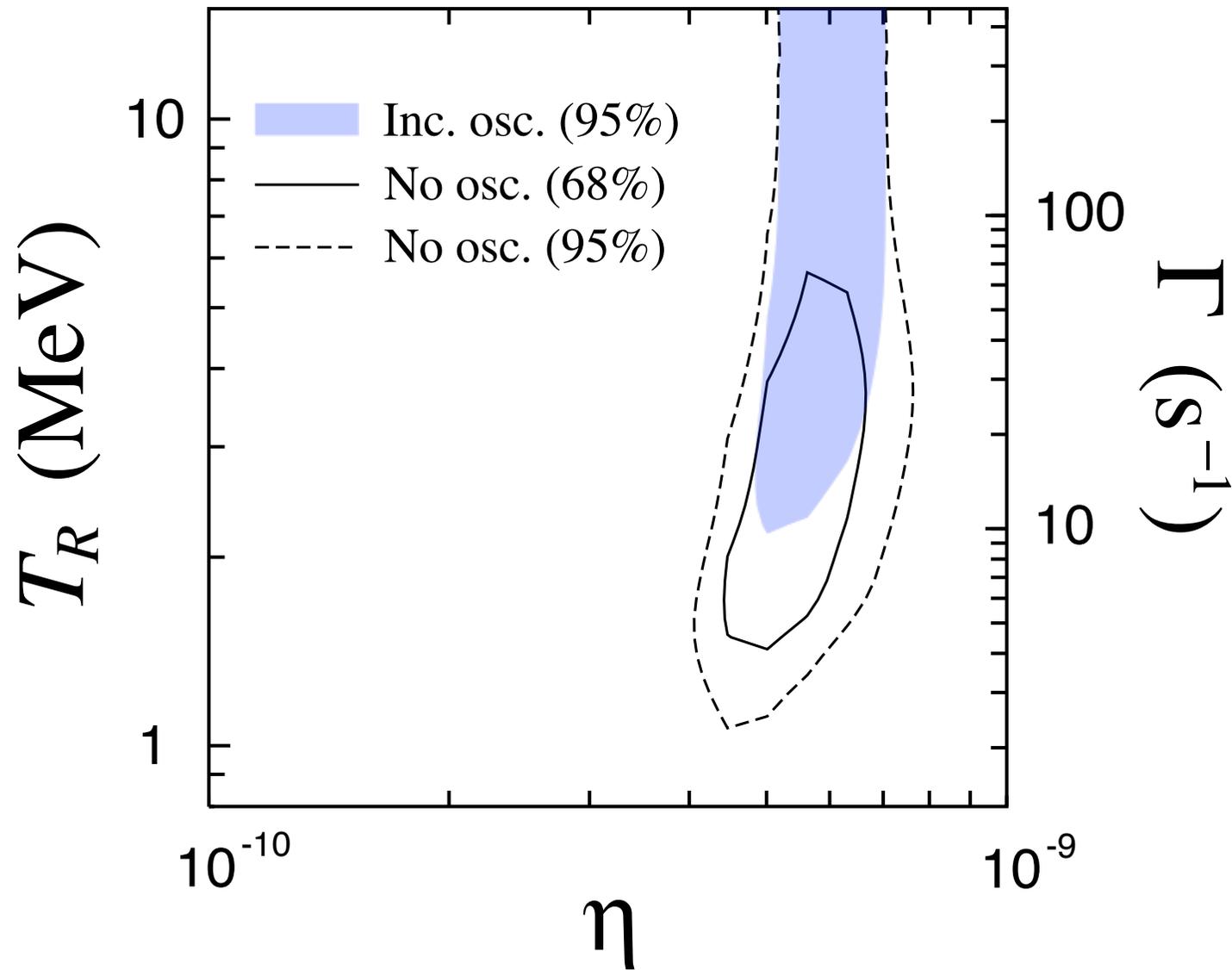


Taking into consideration the neutrino oscillations, the effect (2) becomes more important!

^4He abundance Y_p and T_R



- χ^2 contour plots using data of D and ^4He



V. Conclusion

- We have investigated the **MeV-scale reheating** scenario, paying particular attention to **neutrino oscillations**.
- In contrast to the widespread picture,



while T_R decreases.

- $T_R \gtrsim 2 \text{ MeV}$, or $N_\nu \gtrsim 1.2$ was obtained.