

COSMIC STRING EVOLUTION IN HIGHER DIMENSIONS

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Based on:

hep-ph/0410349 [PRD]

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Summary

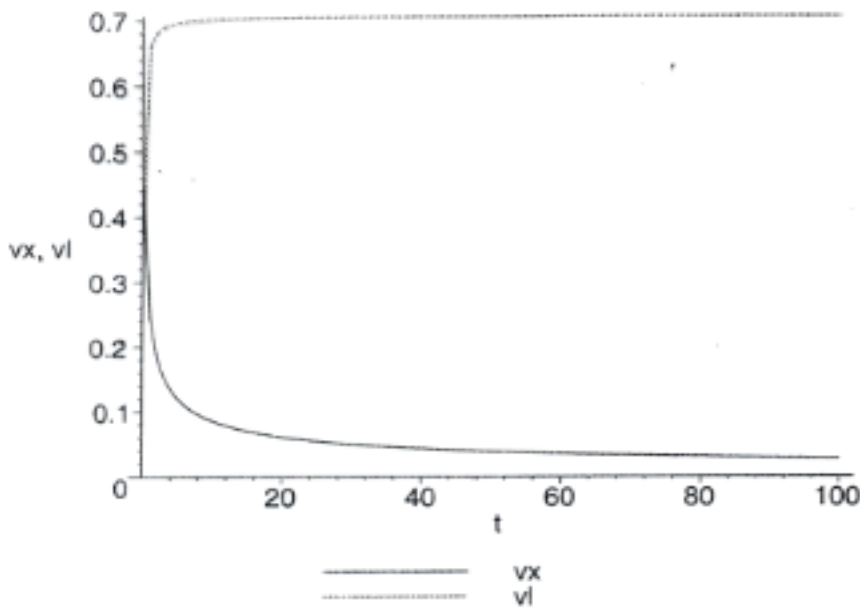
Have developed a model for quantitative string evolution in higher dimensions.

This verifies the existence of a scaling solution with a density enhanced by a factor $\sim \rho_{\text{eff}}^{-2}$

Further, quantitative corrections arise due to variable 3D velocity U_x

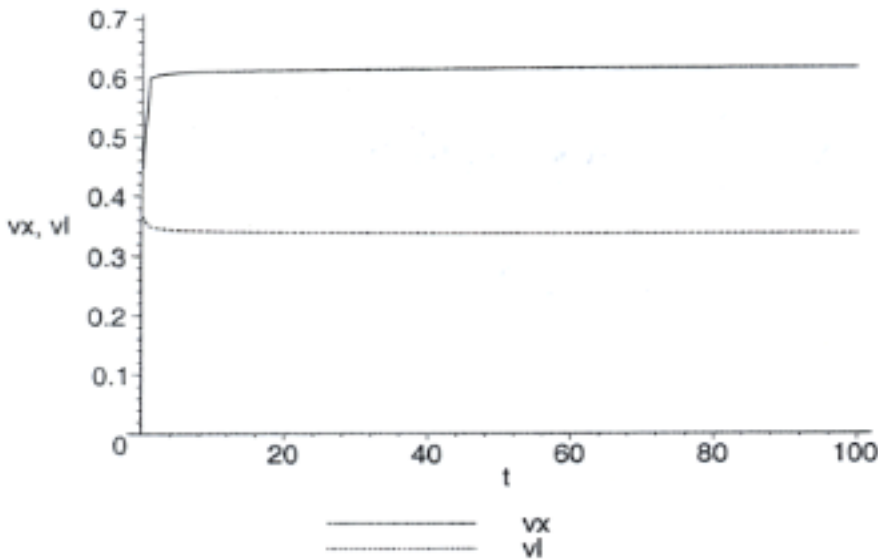
Effect of extra dimensional velocities U_i is to slow down the strings in 3D.

> k1~kx



$U_x \rightarrow 0$
No string motion
in 3D \Rightarrow
No interactions!

> k1 << kx



U_x slowed down

| >

The EDVOS model 2. Macroscopic eqns

• Correlation length

$$E(t) \Rightarrow \dot{E} \Rightarrow \dot{\rho} \Rightarrow \dot{L}, \text{ let } L = \gamma(t) t$$

$$\gamma^{-1} \frac{d\gamma}{dt} = \frac{1}{2t} \left\{ 6[(2+W_1^2) + (2-W_1^2)U_x^2 + (1-W_1^2)U_1^2 - 2 + \frac{\tilde{c}U_x}{\gamma}] \right\}$$

$$\text{where } W_1^2 = \left\langle \frac{\dot{\rho}^2}{a^2 \dot{\rho}^2 + \dot{L}^2} \right\rangle$$

• String velocities

$$U_x \equiv \frac{\int \dot{x}^2 \epsilon d\mathcal{J}}{\int \epsilon d\mathcal{J}} \Rightarrow \dot{U}_x$$

$$U_x \frac{dU_x}{dt} = \frac{K_x U_x}{R} (1-U^2) - (2-W_1^2) H U_x^2 (1-U^2) - H U_x^2 U_1^2$$

$$U_1 \frac{dU_1}{dt} = \frac{K_1 U_1}{R} (1-U^2) - (1-W_1^2) H U_1^2 (1-U^2) + H U_1^2 U_x^2$$

$$\text{where } \frac{K_x U_x (1-U^2)}{R} = \underline{\dot{x}} \cdot \underline{U} \left(1 - \underline{\dot{x}} - \frac{\dot{x}^2}{a^2} \right)$$

sim. K_1

AA & EPSS
hep-ph/0410349
(PRD)

The EDVOS model 1. Microscopic eqns

Nambu action $S = -\mu \int \sqrt{-g} d^2\mathcal{J}$, $g_{ab} = g_{\mu\nu} \partial_a x^\mu \partial_b x^\nu$

Metric: $ds^2 = N(t)^2 dt^2 - a(t)^2 dx^2 - b(t)^2 dl^2$
 x^0 x^1, x^2, x^3 x^0, \dots, x^D

Choose gauge: $\mathcal{J}^0 = t$, $\dot{x} \cdot x' = 0$

• e.o.m. $\frac{\delta S}{\delta x^\mu} = 0$: $\dot{\epsilon} = -N^{-2} \epsilon \left\{ NN' + a\dot{a} \left[\dot{x}^2 - \left(\frac{x'}{\epsilon}\right)^2 \right] + b\dot{b} \left[\dot{l}^2 - \left(\frac{l'}{\epsilon}\right)^2 \right] \right\}$

$\ddot{x} + \left\{ 2\frac{\dot{a}}{a} - N^{-2} \left\{ NN' + a\dot{a} \left[\dot{x}^2 - \left(\frac{x'}{\epsilon}\right)^2 \right] + b\dot{b} \left[\dot{l}^2 - \left(\frac{l'}{\epsilon}\right)^2 \right] \right\} \right\} \dot{x} = \left(\frac{x'}{\epsilon}\right)' \epsilon^{-1}$

$\ddot{l} + \left\{ 2\frac{\dot{b}}{b} - N^{-2} \left\{ NN' + a\dot{a} \left[\dot{x}^2 - \left(\frac{x'}{\epsilon}\right)^2 \right] + b\dot{b} \left[\dot{l}^2 - \left(\frac{l'}{\epsilon}\right)^2 \right] \right\} \right\} \dot{l} = \left(\frac{l'}{\epsilon}\right)' \epsilon^{-1}$

where $\epsilon = \frac{-x'^2}{\sqrt{-g}}$

• E-M tensor from $\frac{\delta S}{\delta g_{\mu\nu}}$:

$T^{\mu\nu} = \frac{1}{Na^3 b^{D-3}} \mu \int d\mathcal{J} (\epsilon \dot{x}^\mu \dot{x}^\nu - \epsilon^{-1} x'^\mu x'^\nu) \delta^{(D)}(x - x(\mathcal{J}, t), l - l(\mathcal{J}, t))$

Energy: $E = \int_{t=\text{const}} \sqrt{h} n_\mu n_\nu T^{\mu\nu} d^3 \underline{x} d^{D-3} \underline{l}$

$\Rightarrow E(t) = N(t) \mu \int \epsilon d\mathcal{J}$

Strings in Higher Dimensions

Motivation from Brane Inflation

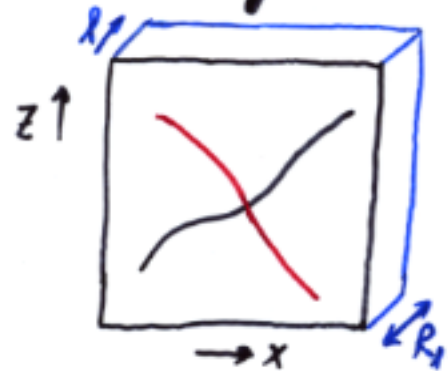
[Dvali & Tye, 1999
BMQRZ, 2001
DSS, 2001]

where string-like defects can be produced at the end of inflation.

(Sarangi & Tye, 2002)
(Dvali & Vilenkin, 2004)

The strings propagate in (FRW) x I space where I a compact manifold.

Consider string collision:



Apparent point of intersection:



strings may miss each other in extra dimension.

Probability of reconnection

$$P \sim \frac{\delta}{R_1}$$

← string width

Introduce P in one-scale model

$$2 \frac{dL}{dt} \approx 2HL + 1 \cdot P \quad (\text{Jones, Stoica & Tye, 2003})$$

Scaling density $\frac{\rho}{\rho_{\text{stand}}} \sim P^{-2}$

Cycloops

[AA & EPSS
hep-ph/0504049
(JHEP)]

Recall string intercommutings produce loops which radiate and shrink to zero radius.



If compact space admits non-trivial 1-cycles then loops can be trapped.



For loops $p(t) \propto t^{-2}$ (like radiation)

For cycloops $p(t) \propto t^{-3/2}$ (like matter)

To avoid cycloop domination for typical choice of parameters in Brane Inflation need $G\mu < 10^{-18}$ too strong

Small regions in parameter space exist for which $G\mu \sim 10^{-14}$ Dark Matter

The VOS model (Martins & Shellard) 1996

Starting from Nambu-Goto action, derive evolution equations for the macroscopic quantities:

$$v = \sqrt{\langle \dot{x}^i \rangle^2}$$

From e.o.m. $\ddot{x} = \dots$

$$L \downarrow$$

From $T^{\mu\nu} \rightarrow E, \dot{E}$
 $\Rightarrow \dot{\rho} \rightarrow \dot{L} \quad [\rho = \mu L^{-2}]$

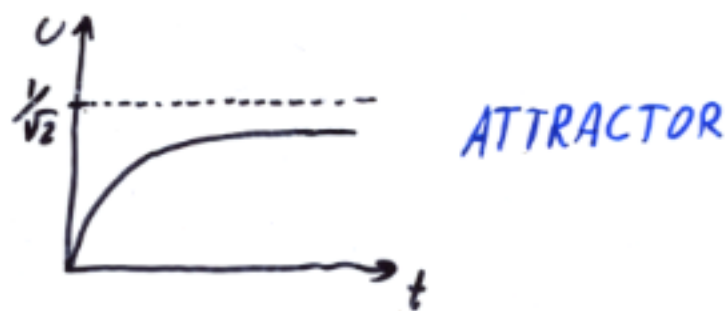
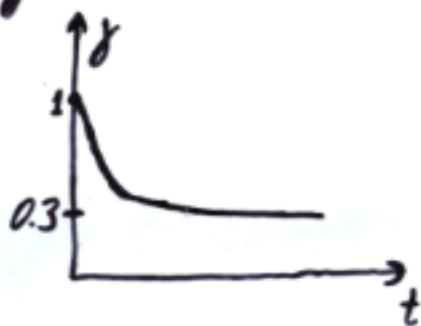
We find:

$$\gamma^{-1} \frac{d\gamma}{dt} = \frac{1}{2t} \left[2b(1+v^2) - 2 + \frac{\tilde{c}v}{\gamma} \right], \quad L = \gamma(t)t$$

$$\frac{dv}{dt} = (1-v^2) \left(\frac{\kappa}{R} - 2Hv \right)$$

where κ - momentum parameter
 b - expansion exponent
 $a(t) \propto t^b$

Scaling Solution:



Constraint: $v^2 \leq \frac{1}{2}$, saturated in flat space

Q. In higher dimensions $v_x^2 + v_z^2 \leq \frac{1}{2}$ $v_x^2 \rightarrow \frac{1}{2}$?
 $v_z^2 \rightarrow 0$?

Cosmic Strings & Evolution

Cosmic Strings: Line-like concentrations of energy arising as topological defects in cosmological phase transitions.



Evolution

- Field Theory Simulations
- Nambu Strings
 - ◀ simulations
 - ◀ 'thermodynamic' approach

String network is Brownian.

Characterised by correlation length L

Energy density $\rho = \frac{\mu}{L^2}$



$\dot{\rho} \approx -2 \frac{\dot{a}}{a} \rho - \frac{\rho}{L}$

Hubble expansion



Find Scaling solution

$L \sim t \sim H^{-1}$
 $\rho \propto t^{-2}$ (Kibble 1985)

Plan

- Cosmic strings & evolution in 3D
- Strings in higher dimensions
- The EDVOS model
- Results & Summary