Gravitino LSP in the Constrained MSSM

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Contents of present Unverse

WMAP data with other CMB experiments (ACBAR and CBI), 2dFGRAS measurements, and Ly α forest data determines the best-fit cosmological parameters[Spergel et.al, (2003)]:

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\Omega_{tot} = \Omega_{\Lambda} + \Omega_{m} + \dots = 1.02 \pm 0.02 \text{ (Flat Universe)}
\begin{cases}
\text{energy} & \Omega_{\Lambda} = 0.73 \pm 0.04 \text{ (accelerating)} \\
\text{matter} & \Omega_{m} = 0.27 \pm 0.04
\end{cases}
\begin{cases}
\Omega_{b}h^{2} = 0.0224 \pm 0.0009 \text{ (BBN, CMB)} \\
\Omega_{DM}h^{2} = 0.113^{+0.008}_{-0.009} \\
\Omega_{\nu}h^{2} < 0.0076 & (h = 0.71^{+0.04}_{-0.03})
\end{cases}
\text{radiation} & \Omega_{\gamma} = (2.471 \pm 0.004) \times 10^{-5} \text{ (}T_{0} = 2.275 \pm 0.002\text{K }95\% \text{ CL)}
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Dark Matter?

What is Dark Matter?

neutrino, neutralino, generic WIMP, axion, axino, Gravitino, WIMPzilla,...

LSP (Lightest Supersymmetric Particle) is a possible candidate for the cold dark matter if R-parity is conserved. The lightest Neutralino is the most promising candidate with an abundance calculated from the freeze-out of annihilation processes in a thermal initial state.

Gravitino LSP : spin 3/2, the superpartner of Graviton, always exists in local SUSY, with mass $m_{\tilde{G}}=\frac{F}{\sqrt{3}M_p}$ F: SUSY breaking scale, M_p : reduced Planck mass

Gravitino Dark Matter

- Gravitino problem
- Is Gravitino stable? LSP or not?
- Relic density of Gravitino for Dark Matter, $\Omega_{\tilde{G}} = \Omega_{DM}$?
- NLSP decay to Gravitino at late times
 - non-thermal production (NTP) of Gravitino
 - Big Bang Nucleosynthesis (BBN) constraints
 - CMB constraints
- Thermal production (TP) from scattering during reheating
- Unbounded From Below (UFB) constraints

Gravitino Problem

If $T_R > T_f$, Gravitino freezes out from thermal equilibrium with the relic density $\Omega_{\tilde{G}} h^2 = 1.17 \left(\frac{100}{g_*} \right) \left(\frac{m_{\tilde{G}}}{1 \text{ keV}} \right)$

$$\Omega_{\tilde{G}} \equiv \frac{\rho_{\tilde{G}}}{\rho_{C}}$$

- Stable : $m_{\tilde{G}} \lesssim 1$ keV from $\Omega_{tot} < 1$ or Gravitino dilution
- ullet Unstable : $m_{ ilde{G}}\gtrsim 10$ TeV from BBN constraint

If $T_R < T_f$, the freeze-out Gravitino is diluted away by Inflation and reproduced by thermal scattering in the reheating period with abundance

$$Y_{\tilde{G}} \equiv \frac{n_{\tilde{G}}}{n_{\gamma}} = 7.7 \times 10^{-12} \left(1 + \frac{m_{\tilde{g}}^2}{12m_{\tilde{G}}^2} \right) \left(\frac{T_R}{10^{10} \text{GeV}} \right)$$

 T_R : Reheating temperature

[Bolz, Brandenburg, Buchmüller, (2001)]

ullet unstable : $T_R \lesssim 10^8~{
m GeV}$ from BBN

[Ellis,Kim,Nanopoulos('84), Cyburt et al.('02), Kawasaki et al.('04)]

• stable : $T_R \lesssim 10^{10}$ GeV from $\Omega_{tot} < 1$

$$\Omega h^2 = 3.63 \times 10^9 \frac{m}{100 \text{ GeV}} Y$$

Gravitino LSP in the CMSSM

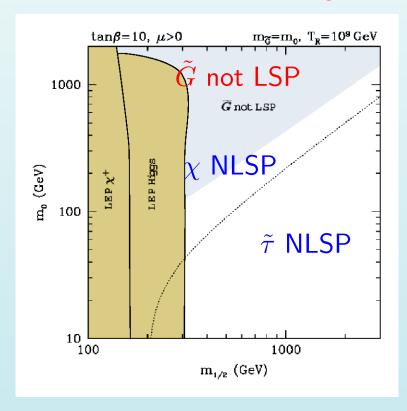
Constrained MSSM

- At M_{GUT} gauginos $M_1=M_2=M_3=m_{1/2}$
- scalars $m_{\tilde{q}}^2=m_{\tilde{l}}^2=m_{H_b}^2=m_{H_t}^2=m_0^2$ • trilinear soft terms $A_b=A_t=A_0$
- radiative EWSB
- five independent parameters: $\tan\beta, m_{1/2}, m_0, A_0, sgn(\mu)$

Free parameter: $m_{\tilde{G}} \sim M_{SUSY}$ Experimental constrains

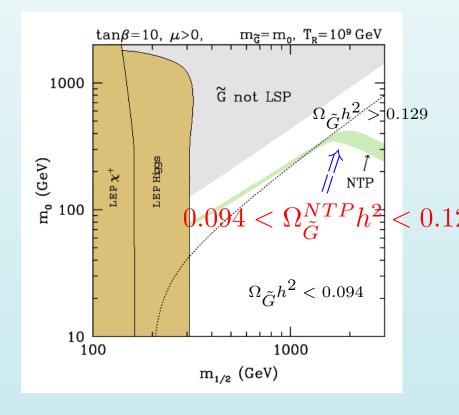
- $m_{_Y\pm}>104~{\rm GeV(LEP)}$
- light higgs: $m_h > 114.4 \text{ GeV(LEP)}$
- BR(B $\to X_s \gamma$)= $(3.34 \pm 0.68) \times 10^{-4}$

e.g.,
$$\tan \beta = 10, \ \mu > 0, \ m_{\tilde{G}} = m_0$$



Gravitino Relic abundance

$$m_{\tilde{G}} = m_0$$



Non-thermal production: NLSP decay

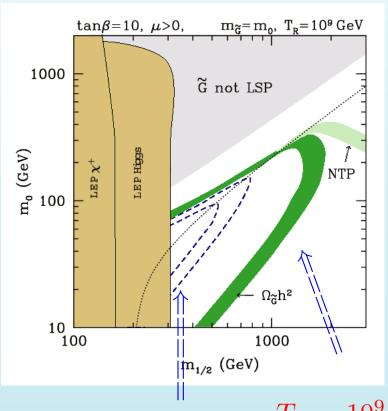
$$\begin{array}{c} \chi \text{ NLSP} \\ \chi \to \tilde{G} \gamma \\ \chi \to \tilde{G} Z, \tilde{G} Higgs, \tilde{G} (\gamma^*/Z^*) q \bar{q} \end{array}$$

$$\begin{array}{ccc} & \tilde{\tau} & \mathsf{NLSP} \\ < 0.129 & \tilde{\tau} \to \tilde{G}\tau \\ & \tilde{\tau} \to \tilde{G}\tau Z, \tilde{G}\nu_{\tau}Ws, \tilde{G}\tau(\gamma^*/Z^*)q\bar{q} \end{array}$$

$$\Omega_{\tilde{G}}^{NTP}h^2 = \frac{m_{\tilde{G}}}{m_{NLSP}}\Omega_{NLSP}h^2$$

Gravitino Relic abundance

$$m_{\tilde{G}} = m_0$$



Thermal production: scattering during reheating

$$\Omega_{\tilde{G}}^{TP}h^2 \simeq 0.2 \left(\frac{T_R}{10^{10}~\text{GeV}}\right) \left(\frac{100~\text{GeV}}{m_{\tilde{G}}}\right) \left(\frac{m_{\tilde{g}}(\mu)}{1~\text{TeV}}\right)^2$$
 [Bolz,Brandenburg,Buchmüller., (2001)]
$$\downarrow \hspace{1cm}$$

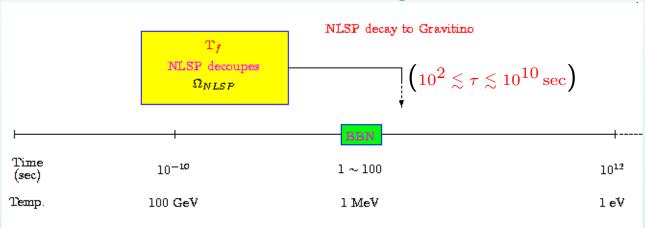
$$\Omega_{\tilde{G}}h^2 = \Omega_{\tilde{G}}^{TP}h^2 + \Omega_{\tilde{G}}^{NTP}h^2$$

$$0.094 < \Omega_{\tilde{G}}h^2 < 0.129$$

$$T_R = 10^9 \text{ GeV}$$

$$T_R = 5 \times 10^9 \text{ GeV}$$

NLSP decay



NLSP decay produces photons and hadrons with high energy

$$\begin{array}{c} \chi \text{ NLSP} \\ \chi \to \tilde{G}\gamma \Rightarrow \text{ EM showers} \\ \chi \to \tilde{G}Z, \tilde{G}Higgs, \tilde{G}\gamma^*/Z^* \\ \Rightarrow \text{ had showers} \\ \end{array} \begin{array}{c} \tilde{\tau} \text{ NLSP} \\ \tilde{\tau} \to \tilde{G}\tau \Rightarrow \text{ EM showers} \\ \tilde{\tau} \to \tilde{G}\tau Z, \tilde{G}\nu_{\tau}W, \tilde{G}\tau\gamma^*/Z^* \\ \Rightarrow \text{ had showers} \\ \end{array}$$

- Late time decay due to gravitational interaction: $10^2 \sim 10^{10}\,\mathrm{sec}$
- The high energy injection(EM,had) at late times during or after BBN changes the light element abundances \Rightarrow Strong constraints on NLSP

BBN constraints

BBN constraints

conservative observational limit

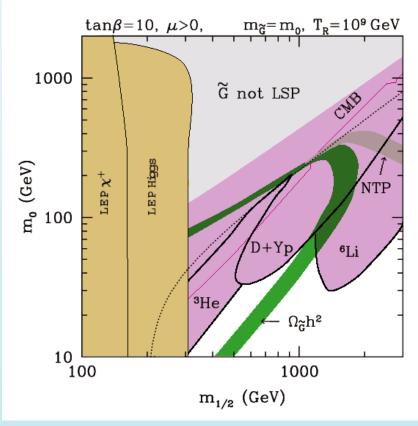
$$2.2 \times 10^{-5} < D/H < 5.3 \times 10^{-5}$$
 $0.232 < Y_p < 0.258$
 $1.11 \times 10^{-10} < ^7\text{Li/H} < 4.5 \times 10^{-10}$
 $^3\text{He/D} < 1.72$
 $^6\text{Li/}^7\text{Li} < 0.1875$

CMB constraints

dimensionless chemical potential μ

$$|\mu| < 9 \times 10^{-5}$$

$$m_{\tilde{G}} = m_0, \ T_R = 10^9 \text{ GeV}$$



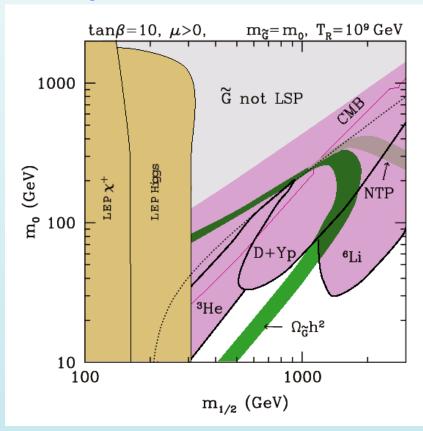
$$D/H+Y_p+^7Li/H+^3He/D+^6Li/^7Li$$

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BBN constraints

- χ NLSP is inconsistent with BBN : still possibility for sub-GeV Gravitino $(\tau_\chi < 1\,{\rm sec})$
- $\tilde{\tau}$ NLSP region is allowed
- CMB constraint is weaker than BBN: potentially important for more precise CMB measurements
- NTP only region is disallowed by BBN: we need another mechanism for Gravitino production
- How high T_R is possible?

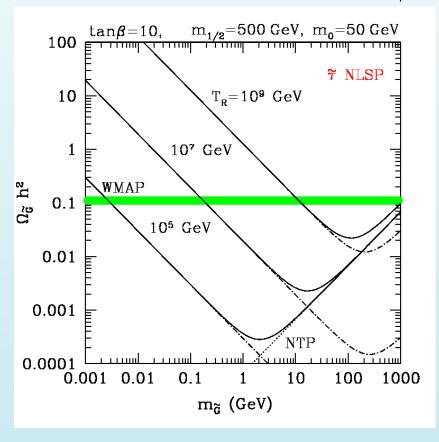
$$m_{\tilde{G}} = m_0, \ T_R = 10^9 \text{ GeV}$$



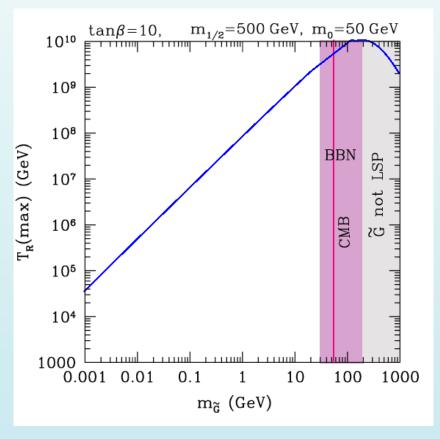
$$D/H+Y_p+^7Li/H+^3He/D+^6Li/^7Li$$

Thermal production

$$\tan \beta = 10, \ m_{1/2} = 500 \text{GeV}, \ m_0 = 50 \text{GeV}(\tilde{\tau} \text{NLSP})$$



$$\Omega_{\tilde{G}}h^2 = \Omega_{\tilde{G}}^{TP}h^2 + \Omega_{\tilde{G}}^{NTP}h^2$$



$$T_R \lesssim 3 imes 10^9 \; {
m GeV} \; {
m for} \; m_{\tilde{G}} \sim 30 \; {
m GeV}$$

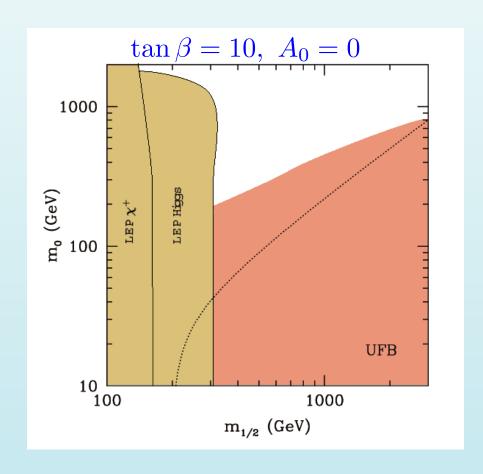
UFB constraints

Two types of constraints:

- charge and/or color breaking (CCB) minima
- unbounded from below (UFB) directions (UFB-1,2,3)
 - \Rightarrow Among them, UFB-3= $\{H_u, \nu_{L_i}, e_{L_j}, e_{R_j}\}$, $i \neq j$ direction leads to electric charge breaking also (the strongest constraints) [Casas, Lleyda, Muñoz('96)]

Condition

$$V_{\text{UFB-3}}(Q = \hat{Q}) > V_{\text{real min}}$$



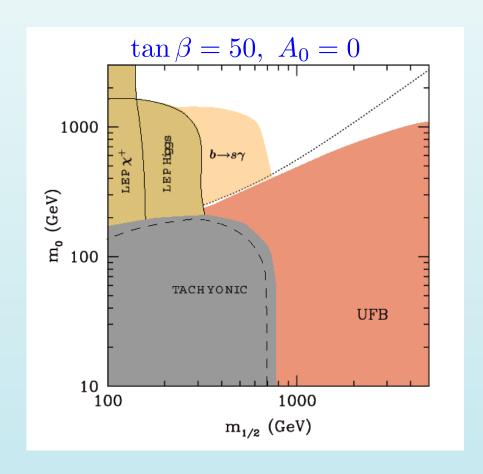
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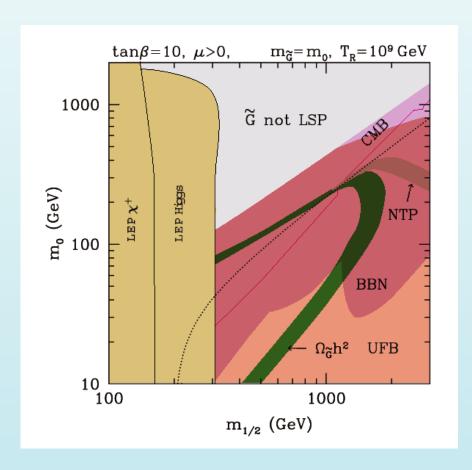
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Condition

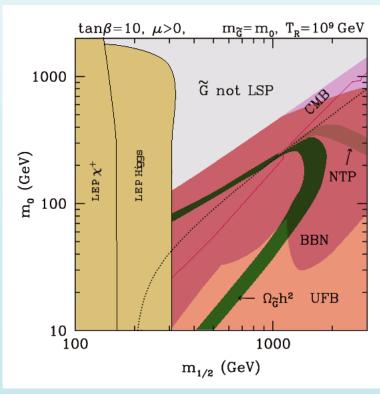
$$V_{\text{UFB-3}}(Q = \hat{Q}) > V_{\text{real min}}$$



Combining UFB and BBN

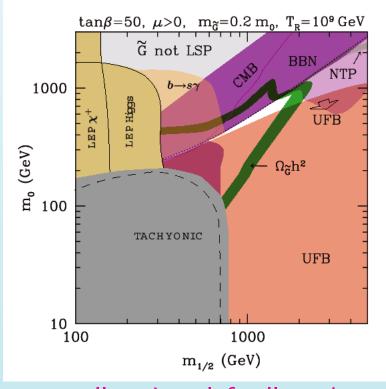
For $T_R = 10^9$ GeV,

• $m_{\tilde{G}} = m_0$, $\tan \beta = 10$



all excluded

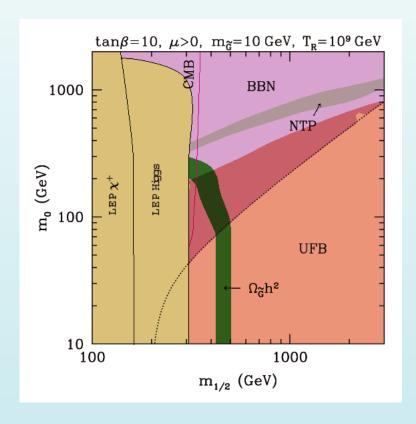
• $m_{\tilde{G}} = 0.2m_0$, $\tan \beta = 50$



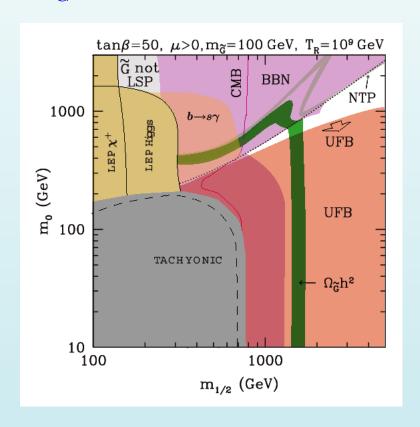
small regions left allowed

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• $m_{\tilde{G}} = 10 \text{ GeV}$, $\tan \beta = 10$



• $m_{\tilde{G}} = 100 \text{ GeV}, \tan \beta = 50$



Discussion

- Neutralino NLSP and O(100)GeV Gravitino LSP is disfavored by BBN: neutralino + sub-GeV Gravitino?
- Stau NLSP and Gravitino LSP is a possible combination in CMSSM for O(100) GeV Gravitino with additional Gravitino production: E.g., for thermal production, up to $T_R \simeq 3 \times 10^9 \; {\rm GeV}$ is available
- UFB constraints may be consistent if the lifetime of realistic vacuum is longer than the age of the Universe due to non-trivial cosmology
- Warm property of NTP Gravitino may solve the small scale problem of CDM or [Borgani et al. (1996)]
- Charged stau NLSP decay may suppress the matter power spectrum on small scales due to the coupling of charged stau to the photon-baryon system [Sigurdson, Kamionkowski (2004)] and/or induces the scale dependent spectral index [Profumo et al. (2005)].