

Superconducting Cosmic Strings in Supergravity

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1 Introduction

Strings, fermionic zero modes & vortons

2 D-term hybrid inflation in global SUSY

One of zero modes is goldstino

3 D-term hybrid inflation in SUGRA

Zero mode absent due to super Higgs effect

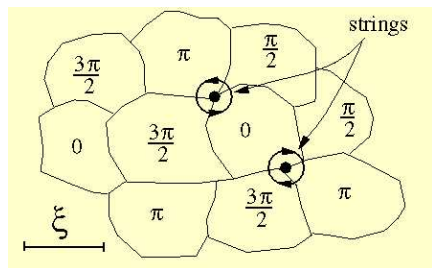
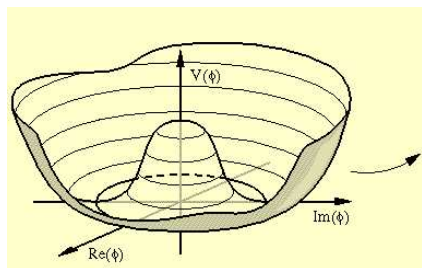
4 Conclusions

R. JEANNEROT & MP ('04)

BECKER, BECKER, STROMINGER ('95)

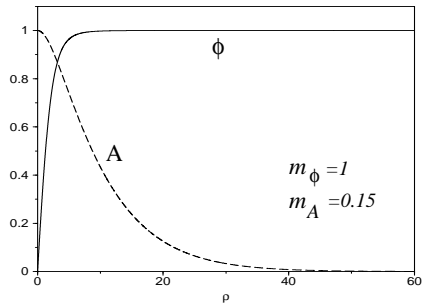
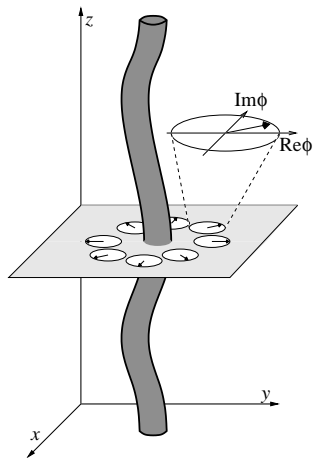
Kibble mechanism

Phase transition in the early universe

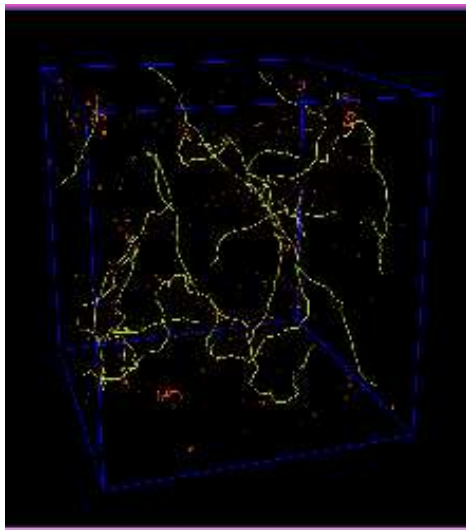


$$l_H = H^{-1} \text{ with } H = \dot{a}/a = \rho/3m_p^2$$

Cosmic strings



String network

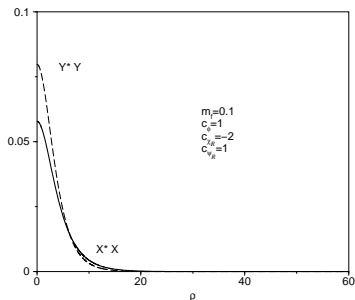


ALLEN, SHELLARD

Fermionic zero modes

$$\mathcal{L}_F = i\psi_i\sigma^\mu D_\mu\bar{\psi}_i - \lambda\phi\psi_i\psi_j + \text{h.c.}$$

$m_\psi = \lambda\phi$ with $\phi = 0$ in string core, $\phi = v$ outside



zero mode solution:

$$\psi_j^{\text{zm}} = \xi_j(r, \theta)e^{\pm ik(t-z)}$$

JACKIW & ROSSI

$$\mathcal{L}_y = -\lambda\phi\psi_i\psi_j - \lambda\phi^\dagger\chi_i\chi_j + \text{h.c.} \quad \text{with Higgs } \phi = vf(r)e^{in\theta}$$

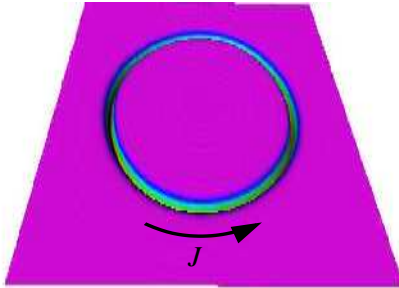
Index theorem

There are n left-moving zero modes for each coupling to ϕ , and n right-moving zero modes for each coupling to ϕ^\dagger .

E. WEINBERG

String loop stabilized by angular momentum of the fermion current

R.L. DAVIS, SHELLARD



- chiral strings
- CDM: constraints

A.C. DAVIS, TRODDEN,
BRANDENBERGER, CARTER

R. JEANNEROT & MP, JHEP 0412:032,2004

R. JEANNEROT & MP, JHEP 0412:043,2004

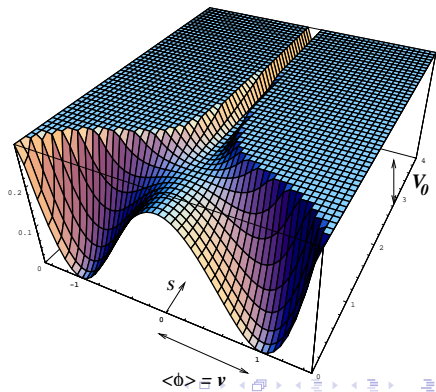
- FI-term $\Leftrightarrow U(1)$
- $W = \kappa S \phi_+ \phi_-$

$$V = \kappa^2 (|S|^2 |\phi_+|^2 + |S|^2 |\phi_-|^2 + |\phi_+|^2 |\phi_-|^2) + \frac{g^2}{2} (|\phi_+|^2 - |\phi_-|^2 - \xi)^2$$

Vacuum

$$|\phi_+| = \sqrt{\xi}$$

$$|S| = |\phi_-| = 0$$



$$V_D = \frac{g^2}{2} (|\phi_+|^2 - \xi)^2$$

string: $\phi_+ = \sqrt{\xi} f(r) e^{in\theta}$

$$A_\theta = -\frac{n}{g} a(r)$$

with $f(0) = a(0) = 0$ and $f(\infty) = a(\infty) = 1$

BOGOMOLNY

$$\begin{aligned} \mu &= \int r dr d\theta \left(|D_\mu \phi_+|^2 + \frac{1}{4} F^2 + \frac{1}{2} D^2 \right) \\ &= \int r dr d\theta \left(|(D_r \mp \frac{i}{r} D_\theta) \phi|^2 + \frac{1}{2} |F_{12} \mp D|^2 \right) \pm 2\pi n \xi \end{aligned}$$

$$V_D = \frac{g^2}{2} (|\phi_+|^2 - \xi)^2$$

string: $\phi_+ = \sqrt{\xi} f(r) e^{in\theta}$

$$f' = n \frac{1-a}{r} f$$

$$A_\theta = -\frac{n}{g} a(r)$$

$$n \frac{a'}{r} = g^2 \xi (1 - f^2)$$

with $f(0) = a(0) = 0$ and $f(\infty) = a(\infty) = 1$

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BPS

$$\delta\psi_+ = \sqrt{2} (\epsilon_- F_+ + i D_\mu \phi_+ \sigma^\mu \bar{\epsilon}_-) = 0$$

$$\delta\lambda = -(iD + \frac{1}{2} \sigma^\mu \bar{\sigma}^\nu F_{\mu\nu}) \epsilon_- = 0$$

with $\sigma^0 \sigma^3 \epsilon_\pm = \pm \epsilon_\pm \Rightarrow \epsilon_- = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

$$\mathcal{L}_Y = -ig\sqrt{2}\phi_+^\dagger\psi_+\lambda - \kappa\phi_+\psi_S\psi_- + \text{h.c.}$$

Index theorem:

- n left-moving zero modes (ψ_+, λ)
- n right-moving zero modes (ψ_S, ψ_-)

Susy trafo with $\sigma^0 \sigma^3 \epsilon_+ = +\epsilon_+ \Rightarrow \epsilon_+ = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

- $[Q, H] = 0$
- satisfy eq. of motion

DAVIS, DAVIS, TRODDEN

$$\begin{aligned}\delta\lambda(r, \theta) &= -i2g\xi(1 - f^2)\epsilon_+ \\ \delta\psi_+(r, \theta) &= i2\sqrt{2}\sqrt{\xi}\frac{n}{r}(1 - a)fe^{i(n-1)\theta}\epsilon_+^*\end{aligned}$$

zero mode:

$$\psi_{\text{zm}}(t, \mathbf{z}, r, \theta) = \alpha(t, \mathbf{z}) \left(\delta\lambda(r, \theta) + \delta\psi_+(r, \theta) \right)$$

$$[\sigma^0 \partial_t - \sigma^3 \partial_z] \alpha(t, \mathbf{z}) = 0 \Rightarrow \alpha(t, \mathbf{z}) = e^{\pm ik(t+z)}$$

zero mode is goldstino

- SUSY broken in string:

$$\phi_+ = \sqrt{\xi} f(r) e^{in\theta}$$

$$A_\theta = -\frac{n}{g} a(r)$$

$$D = g\xi(1 - f(r)^2)$$

$$\text{with } f(0) = 0 \text{ \& } f(\infty) = 1$$

- Goldstino transforms non-linearly:

$$\delta\psi_+ = \sqrt{2}(\epsilon_+ F_+ + iD_\mu\phi_+\sigma^\mu\bar{\epsilon}_+)$$

$$\delta\lambda = -(iD + \frac{1}{2}\sigma^\mu\bar{\sigma}^\nu F_{\mu\nu})\epsilon_+$$

FI-term in SUGRA

$$U_\xi(1) = U_V(1) \oplus U_R(1)$$

BARBIERI, FERRARA, NANOPOULOS, STELLE
STELLE, WEST; FREEDMAN

BINETRUY, DVALI, KALLOSH, VAN PROEYEN

$$K = \sum |\phi_i|^2$$

$$W = \kappa S \phi_+ \phi_i$$

$$f_{\alpha\beta} = 1$$

$$V = V_F + \frac{g^2}{2} (q_+ |\phi_+|^2 + q_- |\phi_-|^2 - \xi)^2$$

vacuum:

$$\phi_- = S = 0, \quad \phi_+ = \xi / \sqrt{q_+}$$

BINÉTRY, DVALI, KALLOSH, VAN PROEYEN

$$V_D = \frac{g^2}{2} (|\phi_+|^2 - \xi)^2$$

string: $\phi_+ = \sqrt{\xi} f(r) e^{in\theta}$

$$A_\theta = -\frac{n}{g} a(r)$$

$$ds^2 = dt^2 - dz^2 - dr^2 - C(r)^2 d\theta^2$$

$$\delta\psi_+ = \sqrt{2}(\epsilon'_- F_+ + iD_\mu \phi_+ \sigma^\mu \bar{\epsilon}'_-) = 0$$

BPS

$$\delta\lambda = -(iD + \frac{1}{2}\sigma^\mu \bar{\sigma}^\nu F_{\mu\nu})\bar{\epsilon}'_- = 0$$

$$\delta\psi_\mu = 2(\partial_\mu + \frac{1}{4}W_\mu^{ab}\sigma_{ab} + \frac{1}{2}iA_\mu^B)\epsilon'_- = 0$$

with $\epsilon'_\pm = e^{\pm\frac{1}{2}i\theta}\epsilon_\pm$ and $\sigma^0\sigma^3\epsilon_\pm = \pm\epsilon_\pm$

BINÉTRY, DVALI, KALLOSH, VAN PROEYEN

$$V_D = \frac{g^2}{2} (|\phi_+|^2 - \xi)^2$$

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$$A_\theta = -\frac{n}{g} a(r)$$

$$ds^2 = dt^2 - dz^2 - dr^2 - C(r)^2 d\theta^2$$

$$f' = n \frac{1-a}{C} f$$

$$n \frac{a'}{C} = g^2 \xi (1 - f^2)$$

$$(1 - C') = A_\theta^B(C, a)$$

$$\delta\psi_+ = \sqrt{2} (\epsilon'_- F_+ + i D_\mu \phi_+ \sigma^\mu \bar{\epsilon}'_-) = 0$$

$$\delta\lambda = -(iD + \frac{1}{2} \sigma^\mu \bar{\sigma}^\nu F_{\mu\nu}) \bar{\epsilon}'_- = 0$$

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with $\epsilon'_\pm = e^{\pm \frac{1}{2} i\theta} \epsilon_\pm$ and $\sigma^0 \sigma^3 \epsilon_\pm = \pm \epsilon_\pm$

BPS

BINÉTRY, DVALI, KALLOSH, VAN PROEYEN

$$V_D = \frac{g^2}{2} (|\phi_+|^2 - \xi)^2$$

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$$ds^2 = dt^2 - dz^2 - dr^2 - C(r)^2 d\theta^2$$

$$f' = n \frac{1-a}{C} f$$

$$n \frac{a'}{C} = g^2 \xi (1 - f^2)$$

$$(1 - C') = A_\theta^B(C, a)$$

Asymptotic solutions:

$$r \rightarrow 0: \quad f = 0, \quad a = 0, \quad D = g\xi, \quad A_\theta^B = 0, \quad C = r$$

$$r \rightarrow \infty: \quad f = \sqrt{\xi}, \quad a = 1, \quad D = 0, \quad A_\theta^B = \frac{n\xi}{m_p^2}, \quad C = r \left(1 - \frac{n\xi}{m_p^2}\right)$$

SUSY trafo with $\epsilon'_+ = e^{+\frac{1}{2}i\theta} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

$$\delta\lambda = -i2g\xi(1 - f^2)\epsilon'_+$$

$$\delta\psi_+ = i2\sqrt{2}\sqrt{\xi}\frac{n}{C}(1 - a)fe^{i(n-1)\theta}\epsilon'^*_{+}$$

$$\delta\psi_\theta = 2iA_\theta^B = \begin{cases} 0 & (r \rightarrow 0) \\ \frac{2in\xi}{m_p^2} & (r \rightarrow \infty) \end{cases}$$

Non-renormalizable!

Super-Higgs effect:

$$\begin{aligned} e^{-1}\mathcal{L}_{\text{mix}} &= \frac{i}{2}D\bar{\psi}_\mu\sigma^{\mu\nu}\psi_\nu + \frac{i}{\sqrt{2}}D\bar{\psi}_\mu\gamma^\mu(\delta\lambda + \delta\psi_+) \\ &= \frac{i}{2}D\bar{\psi}'_\mu\sigma^{\mu\nu}\psi'_\nu \end{aligned}$$

- gravitino not confined —no equivalent of SUSY goldstino zero mode

$$m_{3/2} = D = g\xi \text{ in core and } m_{3/2} = 0 \text{ outside}$$

- n right-moving zero modes (ψ_S, ψ_-)

$$\mathcal{L}_Y = -\kappa\phi_+\psi_S\psi_- + \text{h.c.}$$

$$m = 0 \text{ in core and } m = \kappa\sqrt{\xi} \text{ outside}$$

Cosmic strings break SUSY in their core

global SUSY

- n independent zero modes for each breaking —goldstinos
- possible additional zero modes from superpotential

SUGRA

- goldstino zero modes eaten in super-Higgs effect, absent
- additional zero modes from superpotential unaffected

SUGRA D term inflation: n zero modes

SUGRA F term inflation: no zero modes

⇒ Implications for vorton formation

Confusion about number of zero modes

- S. C. Davis, A. C. Davis and M. Trodden, Phys. Lett. B **405** (1997) 257 ,Phys. Rev. D **57** (1998) 5184.
- B. Carter and P. Peter, Phys. Lett. B **466** (1999) 41
- B. Carter and A. C. Davis, Phys. Rev. D **61** (2000) 123501
- A. C. Davis, T. W. B. Kibble, M. Pickles and D. A. Steer, Phys. Rev. D **62** (2000) 083516
- R. Cordero and E. Rojas, Int. J. Mod. Phys. A **17** (2002) 73
- C. N. Ferreira, M. B. D. Porto and J. A. Helayel-Neto, Nucl. Phys. B **620** (2002) 181
- M. Pickles and A. C. Davis, Phys. Lett. B **520** (2001) 345
- A. Achucarro, A. C. Davis, M. Pickles and J. Urrestilla, Phys. Rev. D **66** (2002) 105013, Phys. Rev. D **68** (2003) 065006
- S. C. Davis, P. Binetruy and A. C. Davis, arXiv:hep-th/0501200.

string: SUSY broken, non BPS

$$W = \kappa \mathcal{S}(\phi_+ \phi_- - v^2)$$

$$\phi_+ = \phi_-^\dagger = e^{in\theta} v f$$

$$A_\theta = -\frac{n}{g} a$$

$$F_0 = \kappa v^2 (1 - f^2)$$

Global SUSY: two goldstinos from SUSY trafo with ϵ_\pm

index theorem:

$$\begin{aligned} \mathcal{L}_Y &= ig\sqrt{2}\phi_+^\dagger \psi_+ \lambda - \kappa \phi_+ \psi_- \psi_S - ig\sqrt{2}\phi_-^\dagger \psi_- \lambda - \kappa \phi_- \psi_+ \psi_S \\ &\sim -\phi \psi_- (\psi_S + \lambda) - \phi^* \psi_+ (\psi_S - \lambda) \end{aligned}$$

SUGRA: no zero modes

Chiral strings possible via Majorana term:

$$W = h\phi_+\chi_{-1/2}^2$$

NB. RH neutrino zero mode decays at EW phase transition

Susy trafo with $\sigma^0 \sigma^3 \epsilon_+ = +\epsilon_+ \Rightarrow \epsilon_+ = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

- $[Q, H] = 0$
- satisfy eq. of motion

DAVIS, DAVIS, TRODDEN

$$\begin{aligned}\delta\lambda(r, \theta) &= -i2g\xi(1-f^2)\epsilon_+ \\ \delta\psi_+(r, \theta) &= i2\sqrt{2}\sqrt{\xi}\frac{n}{r}(1-a)fe^{i(n-1)\theta}\epsilon_+^*\end{aligned}$$

How about $n > 1$?

Lagrangian invariant under trafo with $\zeta(x)$ with $\bar{\sigma}^\mu \partial_\mu \zeta(x) = 0$

since $\delta\mathcal{L} = \partial_\mu \bar{\zeta} j^\mu = \dots (\bar{\sigma} \partial_\mu \zeta)^\dagger \dots = 0$

SUSY trafo with $\zeta_+ = e^{i(n-1)\theta} r^{(n-1)} \epsilon_+$

global SUSY

$$\int d^4x d^4\theta \xi V$$

$$V \rightarrow V + \frac{1}{g}(\Lambda - \bar{\Lambda})$$

FI-term

BARBIERI, FERRARA, NANOPOULOS, STELLE
STELLE, WEST; FREEDMAN

SUGRA

$$\int d^4x d^4E \theta e^{\xi V/m_{\text{p}}^2}$$

$$V \rightarrow V + \frac{1}{g}(\Lambda - \bar{\Lambda})$$

$$E \rightarrow E e^{\Sigma + \bar{\Sigma}}$$

R-symmetry: $[Q, R] = -iQ$

$$\theta \rightarrow e^{-i\alpha} \theta \quad \& \quad \bar{\theta} \rightarrow e^{i\alpha} \bar{\theta}$$

$$V \rightarrow V, \quad \lambda \rightarrow e^{i\alpha} \lambda$$

$$\phi \rightarrow e^{i r \alpha} \phi, \quad \psi \rightarrow e^{i(r-1)\alpha} \psi$$

$$e_{\mu}^m \rightarrow e_{\mu}^m, \quad \psi_{\mu} \rightarrow e^{i\alpha} \psi_{\mu}$$